

# TEM heat transport and fluctuations in the HSX stellarator: experiments and comparison with gyrokinetic simulation

J. Smoniewski<sup>1</sup>, E. Sánchez<sup>2</sup>, B.J. Faber<sup>1</sup>, I. Calvo<sup>2</sup>, M.J. Pueschel<sup>1</sup>, K.M. Likin<sup>1</sup>, C.B. Deng<sup>1</sup>, J.N. Talmadge<sup>1</sup>

<sup>1</sup>HSX Plasma Laboratory, University of Wisconsin-Madison, Madison, Wisconsin 53706, U.S.A.

<sup>2</sup>Laboratorio Nacional de Fusión, CIEMAT, 28040 Madrid, Spain.

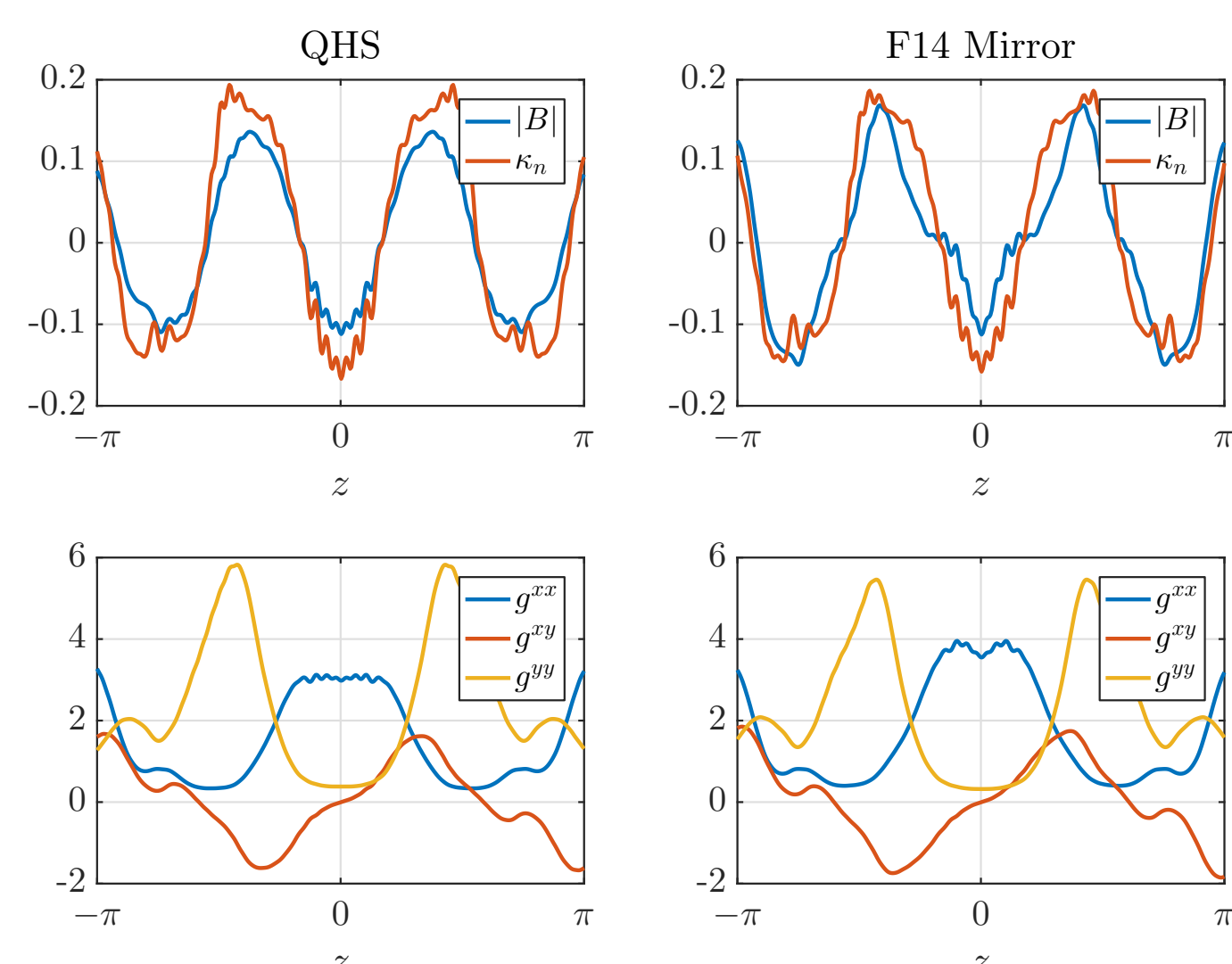


## MOTIVATION

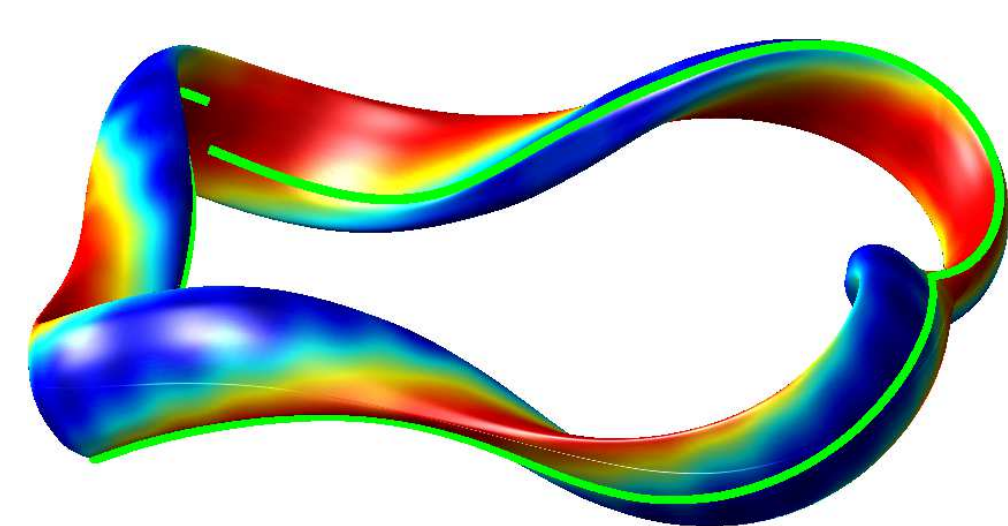
- Optimization for turbulent transport requires predictive capabilities.
- Thus, experimental validation of codes is needed for stellarators.
- Zonal flows regulate transport in many turbulence regimes, also present in quasi-symmetric stellarators.
- Zonal flow residual finite as  $k_x \rho_s \rightarrow 0$  in tokamaks, but vanishes in W7-X (Monreal 2016). Similar for quasi-symmetric stellarators?

## I. Configurations on HSX

- HSX auxiliary coils  $\rightarrow$  comparison of geometries as perturbations on one machine.
- QHS - Quasi-Helical Symmetry, configuration optimized for reduced neoclassical transport and flow damping.
- F14 Mirror - Broken symmetry with  $[n,m] = 4,0$  and 8,0 mirror term, effective ripple similar to conventional stellarator.



- Mirror configuration has narrower  $|B|$  well, no difference in curvature  $\kappa_n$ .
- In Mirror, larger flux compression  $g^{xx}$  at  $z = 0$ , smaller local shear  $g^{yy}$  at  $z = \pm\pi/2$ .
- Gyrokinetics: GENE ([www.genecode.org](http://www.genecode.org)) – local flux tube domain with 4 poloidal turns.



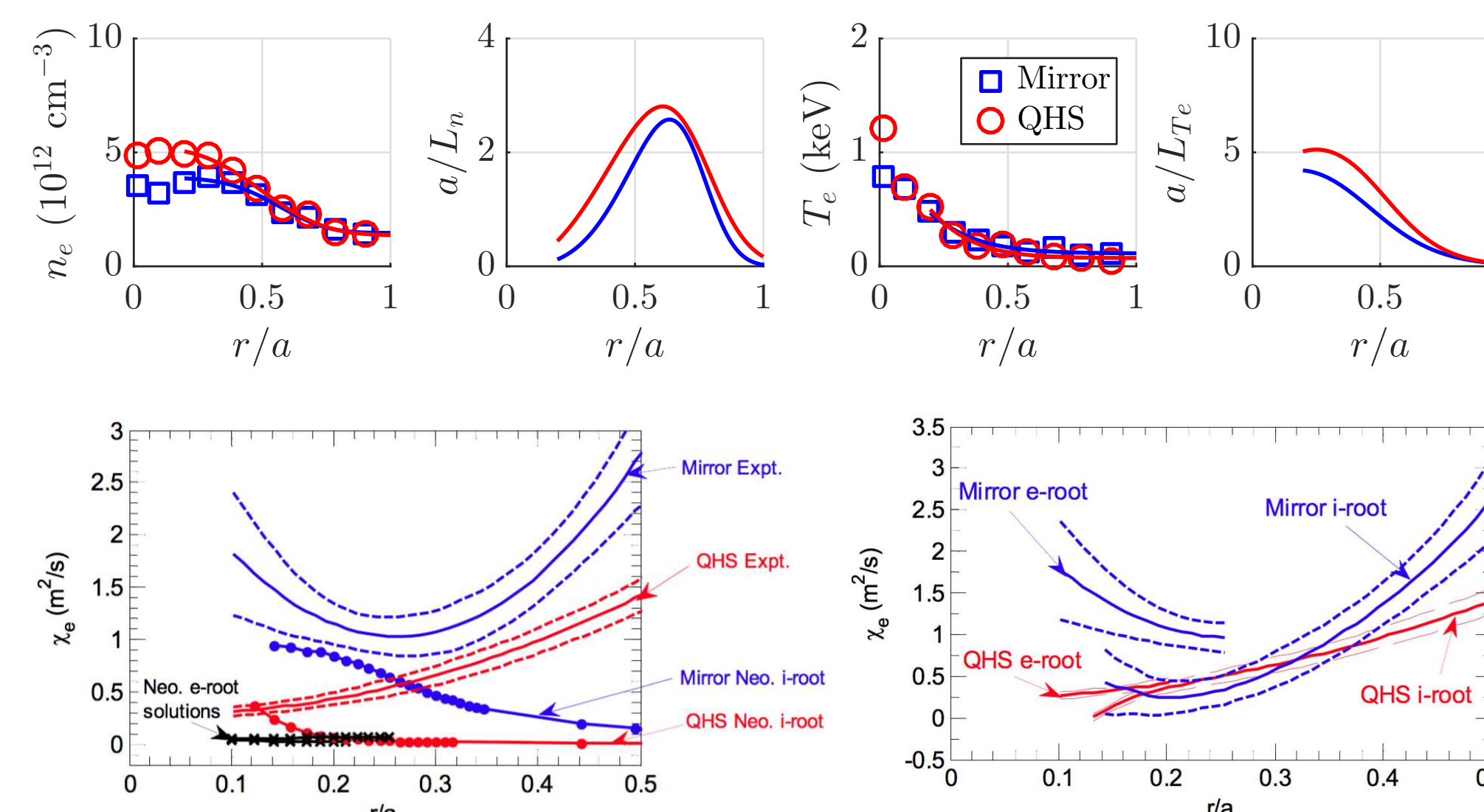
For more on poloidal turns in GENE at HSX, see talk by B.J. Faber GO4.00001 9:30 Tuesday.

## II. Experimental profile analysis: differences in $\chi_e$

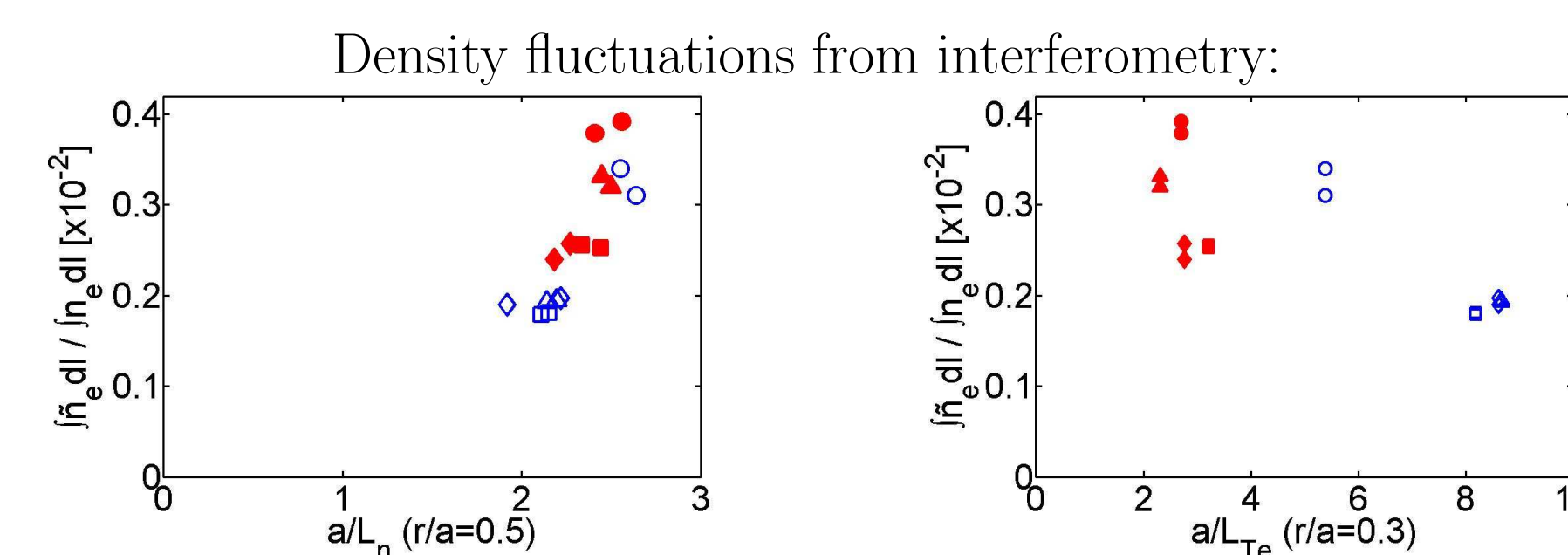
- Plasma profiles are matched in QHS and Mirror with different ECRH power (QHS = 26 kW, Mirror = 44 kW).
- Peak  $\nabla n$  outside  $r/a \approx 0.5$ , peak  $\nabla T_e$  in core – separation of driving gradients.
- Assuming diffusive heat flux ( $Q_e = -n_e \chi_e \nabla T_e$ ),  $\chi_e$  is a normalized measure of energy transport.

$$\chi_e = \frac{P_{abs}}{V' \langle |\rho|^2 \rangle n_e \frac{\partial T_e}{\partial \rho}}$$

- Neoclassical transport approaches zero outside  $r/a \approx 0.3 \rightarrow$  anomalous dominated.
- At  $r/a = 0.5$ ,  $\chi_e$  differs by factor of two between configurations, not accounted by neoclassical transport.
- Interferometry – density fluctuations scale with  $\nabla n$ , but no discernible trend with  $\nabla T_e$ .
- Density scaling consistent with  $\nabla n$ -driven TEM.

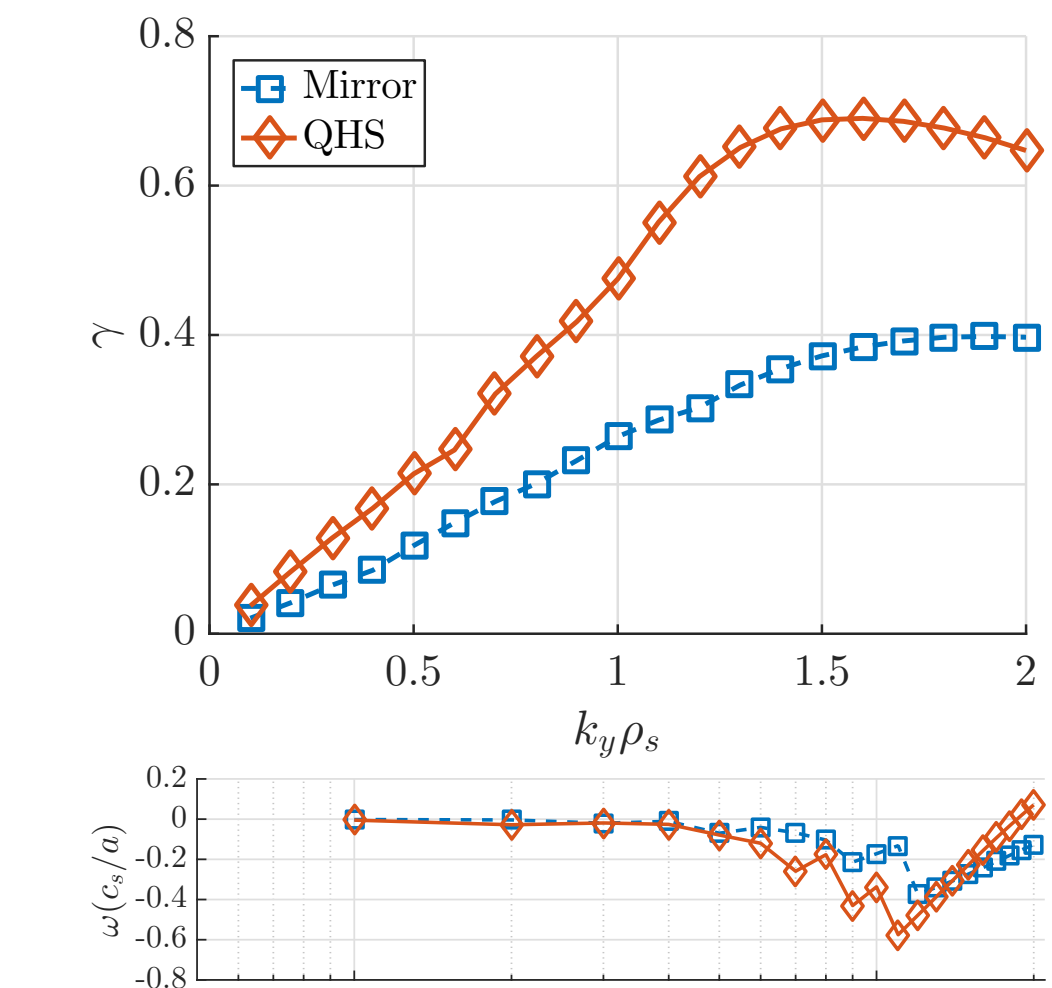


From J. Lore, Ph.D. Thesis, UW-Madison (2010).

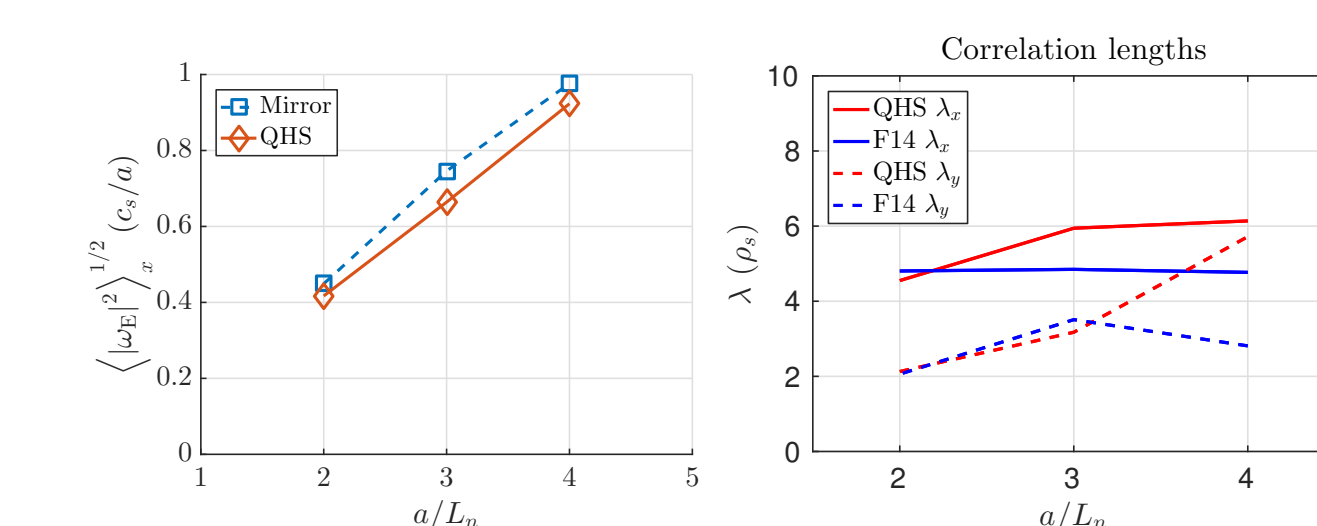
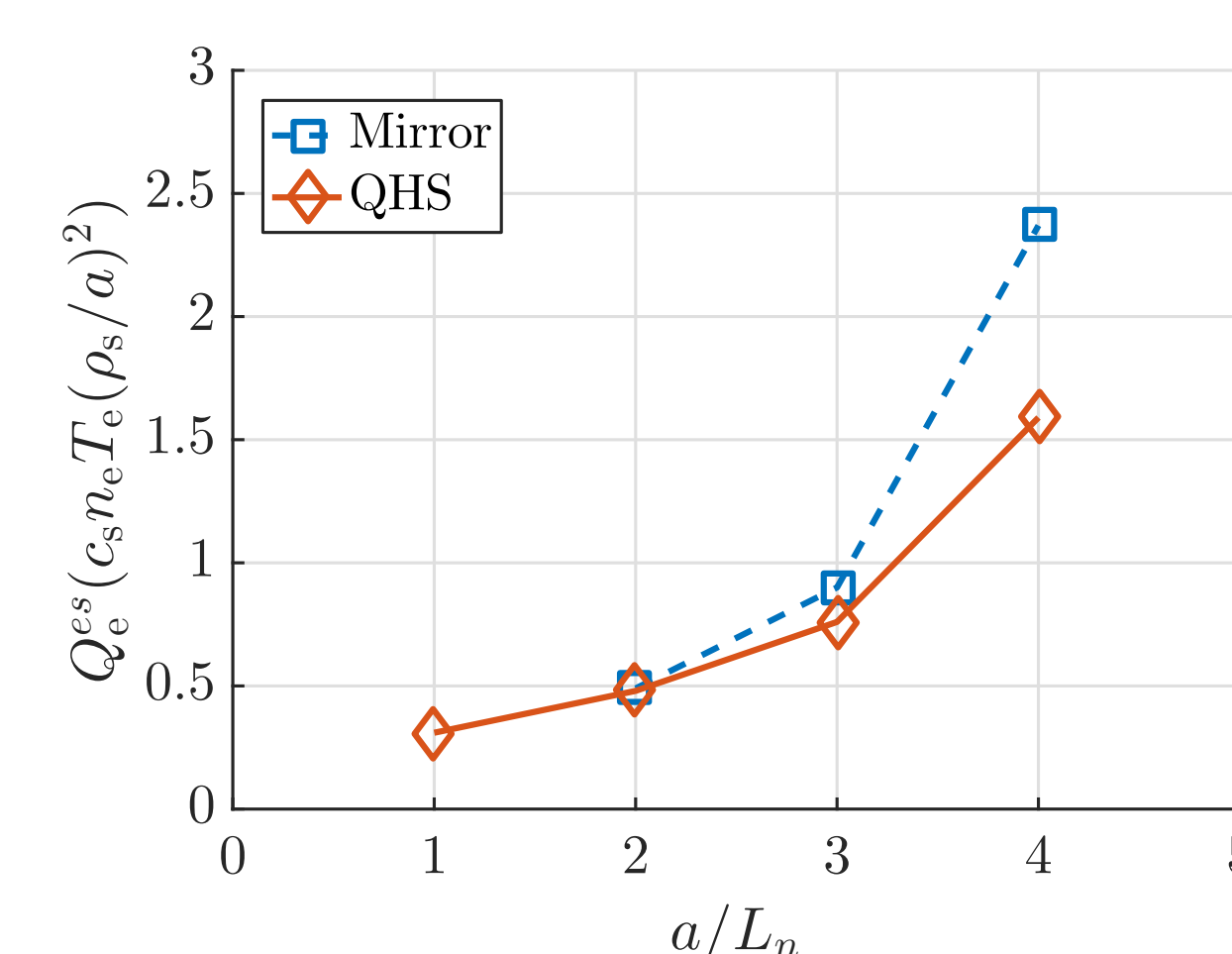
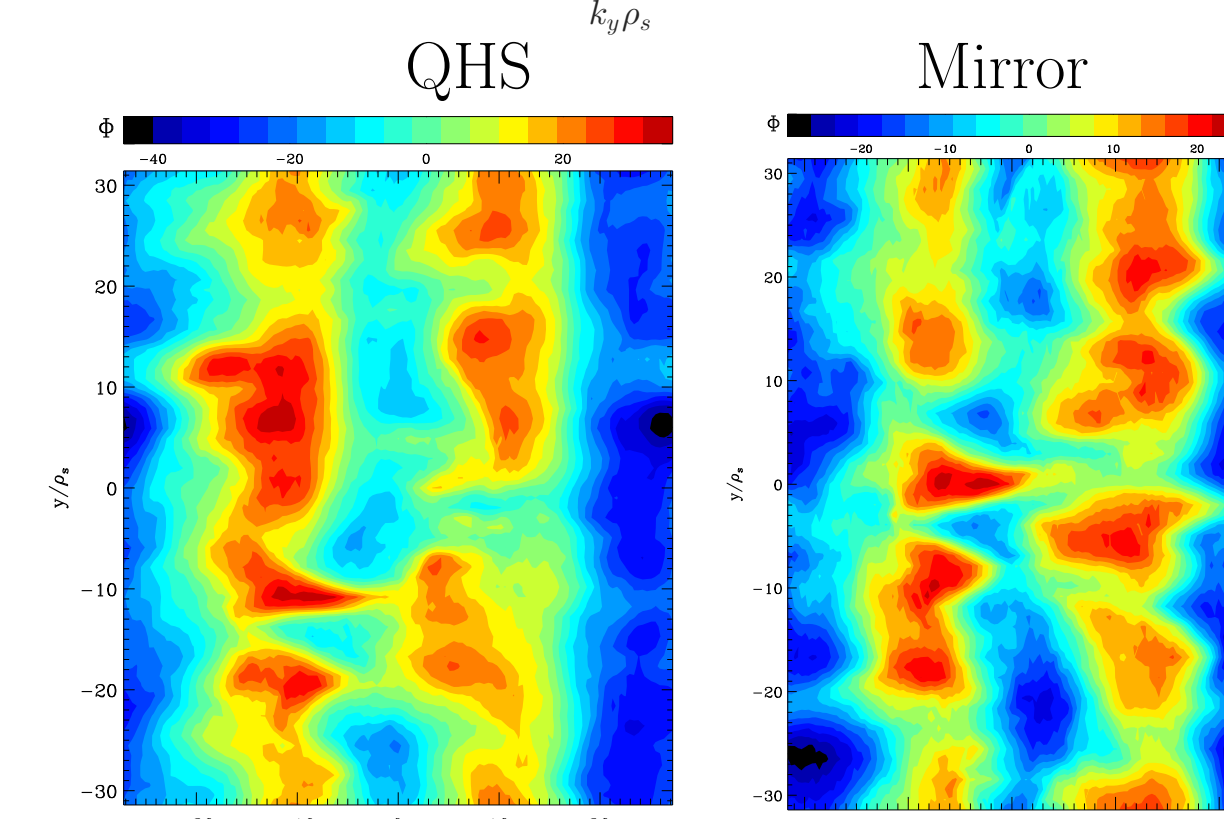
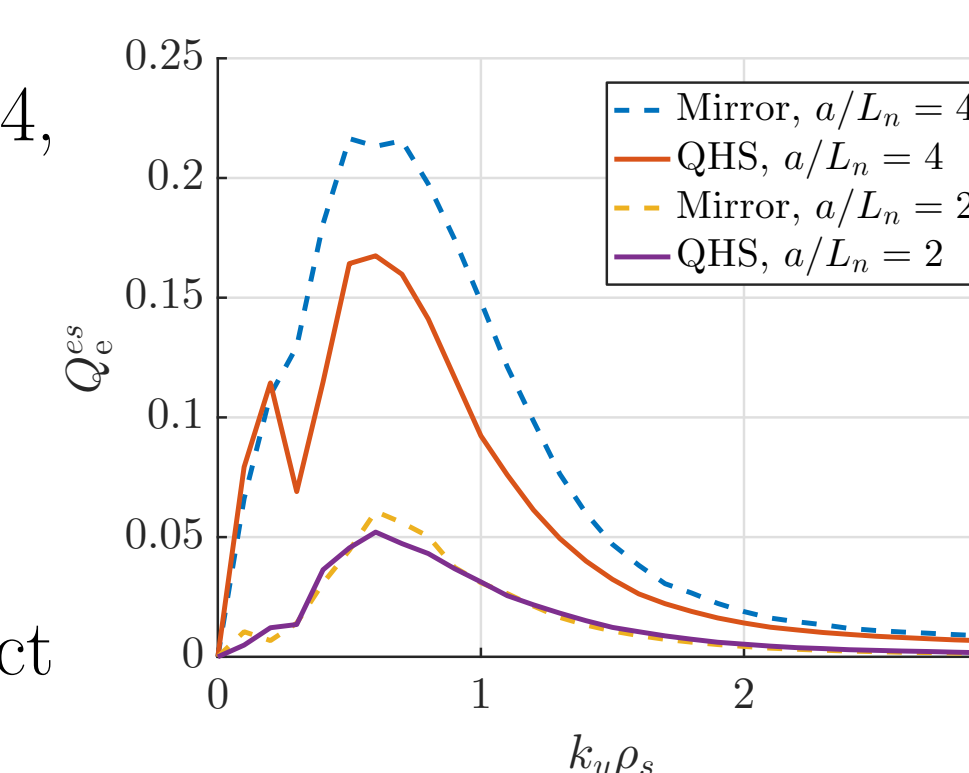


From C.B. Deng et al., Nucl. Fusion. 55 (2015).

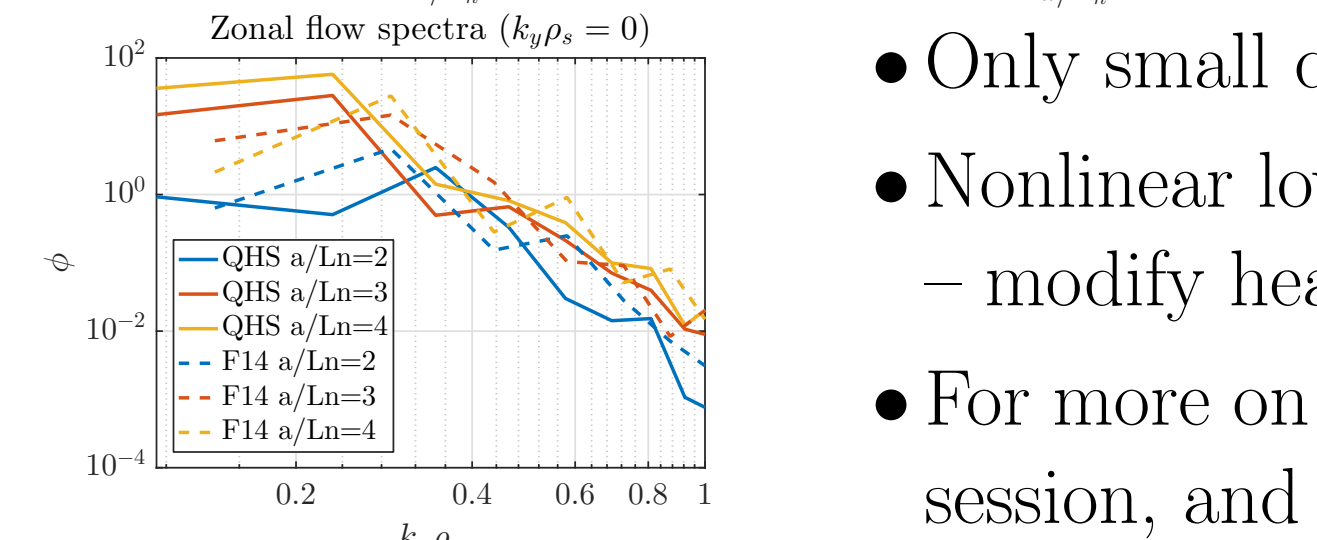
## III. Simulation heat flux in qualitative agreement with experiment



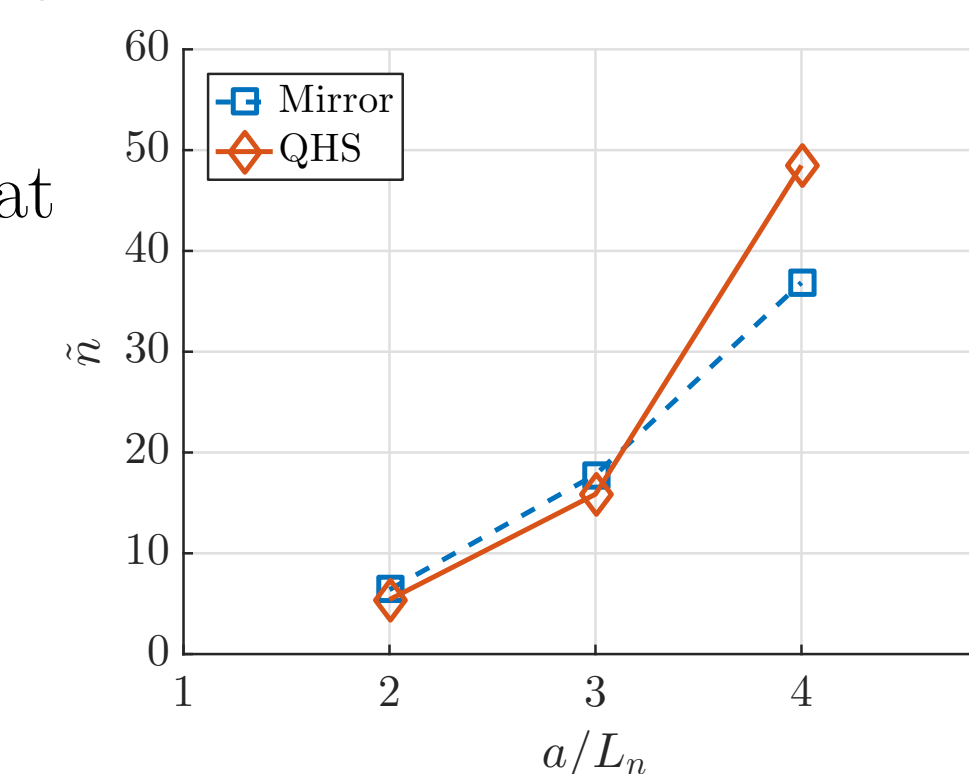
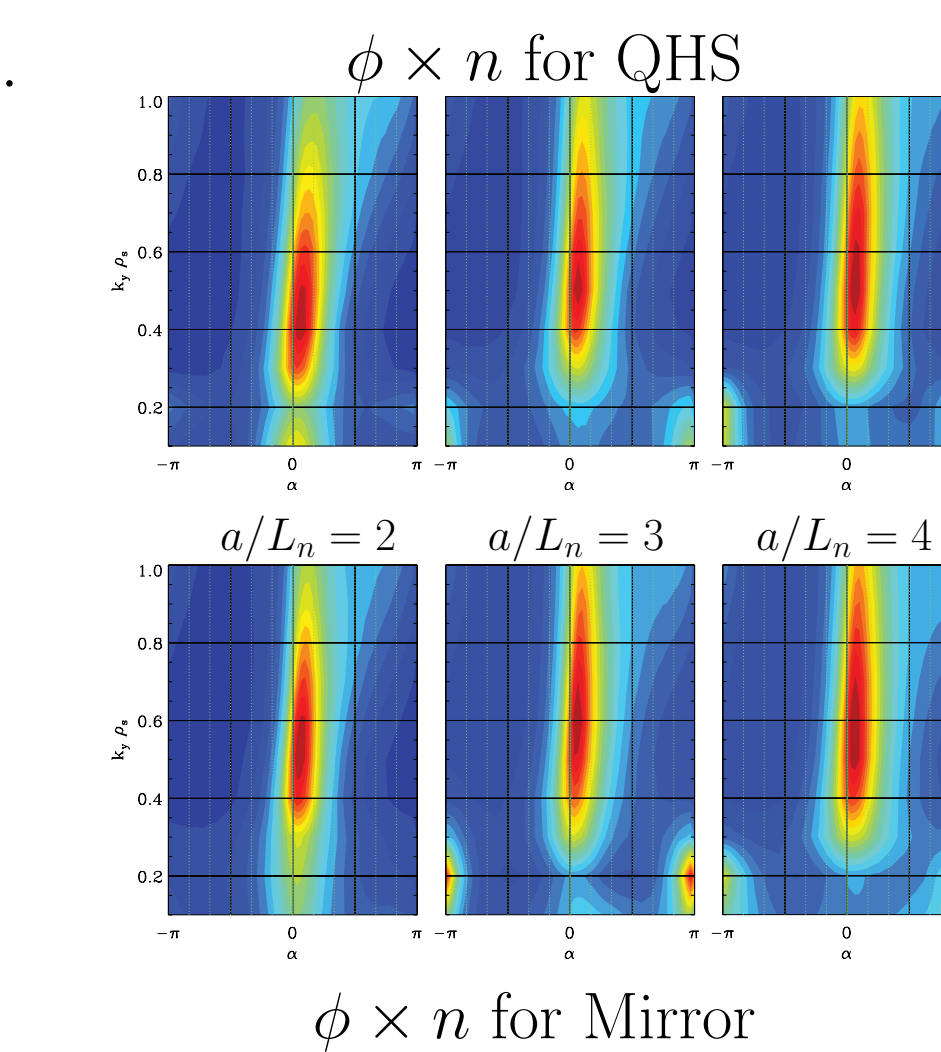
- Most unstable linear growth rates ( $a/L_n = 4$ ,  $a/L_{T_e} = 0$ ) significantly smaller in Mirror.
- Nonlinear turbulence similar in configurations: – consistent spectra, prominent zonal flows, no change in crossphases.
- Nonlinear GENE simulations *not* modeling exact experimental parameters ( $a/L_{T_e} = 0$ ,  $T_i = T_e$ ).
- At high  $a/L_n$ , heat flux is larger in Mirror than in QHS. Encouraging for future validation.



- Radial and poloidal correlation lengths reduced in Mirror at large gradient.
- Density fluctuations larger in QHS at large  $a/L_n$  – opposite heat flux.

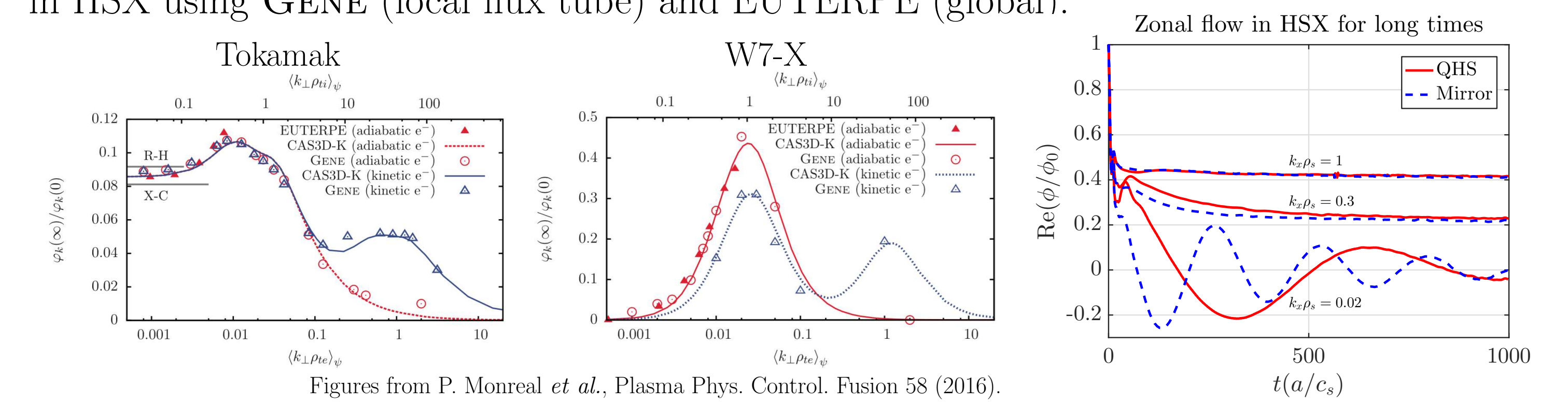


- Only small difference in zonal flow shearing rate.
- Nonlinear low- $k_y \rho_s$  mode develops at large  $a/L_n$  – modify heat transport and saturation properties.
- For more on saturation, see C.C. Hegna CP11.0072 this session, and B.J. Faber GO4.00001 9:30 Tuesday.

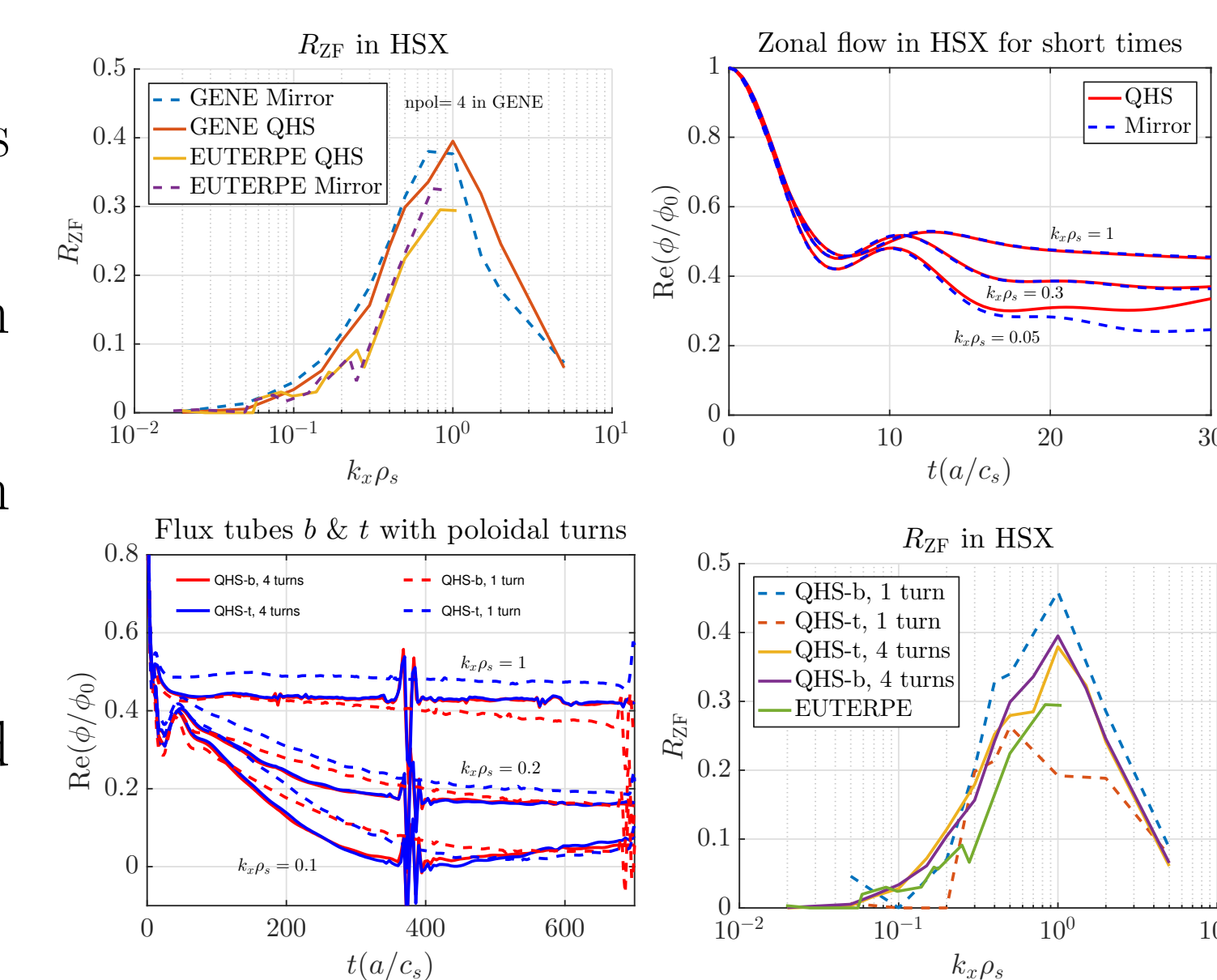


## IV. Zonal flow residual in quasi-symmetry

- In a tokamak, Rosenbluth-Hinton zonal flow residual  $R_{ZF}$  finite as  $k_x \rho_s \rightarrow 0$ .
- In Wendelstein 7-X,  $R_{ZF} \rightarrow 0$  at small  $k_x \rho_s$  and exhibits zonal flow oscillations at  $\Omega_{ZF}$ .
- Do non-ideal quasi-symmetric devices have finite  $R_{ZF}$  as  $k_x \rho_s \rightarrow 0$ ? We compare zonal damping in HSX using GENE (local flux tube) and EUTERPE (global).



- Zonal flow oscillations present at low- $k_x \rho_s$ . Frequency decreases with reduced radial drifts from neoclassical optimization.
- No significant difference in  $R_{ZF}$  between QHS/Mirror. Large difference in  $\Omega_{ZF}$ .
- For times comparable to turbulence correlation time – no difference in QHS/Mirror.
- Damping depends on flux surface averages – long flux tube ( $> 1$  poloidal turn) required to match multiple flux tubes and global.



- Initial results for NCSX similar:  $R_{ZF} \rightarrow 0$  as  $k_x \rho_s \rightarrow 0$ .
- Poloidal turn requirements depend on geometry: in all cases studied here, 1 turn yields qualitatively incorrect results.

## V. Summary

- Anomalous thermal diffusivity  $\chi_e$  at peak density gradient is larger in a broken symmetry configuration than with quasi-helical symmetry.
- Gyrokinetic simulations reproduce a higher electron heat flux with broken symmetry, but not modeling exact experimental parameters. Not predicted by most unstable linear growth rate.
- Zonal flow damping requires flux surface averages – sensitive to computational domain.
- Real, as opposed to ideal, quasi-symmetric devices are more similar to non-symmetric stellarators than tokamaks –  $R_{ZF} \rightarrow 0$  as  $k_x \rho_s \rightarrow 0$  and exhibit zonal flow oscillations.
- Zonal flow damping in QHS, compared to Mirror, exhibits reduced oscillation frequency and decay consistent with reduced radial particle drifts in optimized stellarators.
- Future work: validation of density fluctuations with reflectometer localized to large  $\nabla n$  region with simulations at experimental  $a/L_{T_e}$  and  $T_i$ . Collaboration at W7-X starting in November.

## Acknowledgments

- This work is supported by US DOE Grant DE-FG02-93ER54222, and used resources of the National Energy Research Scientific Computing Center, a DOE Office of Science User Facility supported by the Office of Science of the U.S. Department of Energy under Contract No. DE-AC02-05CH11231.
- This work has been partially funded by the Ministerio de Economía y Competitividad of Spain under project ENE2015-70142-P. The authors thankfully acknowledge the computer resources, technical expertise and assistance provided by the Barcelona Supercomputing Center-Centro Nacional de Supercomputación and the CIEMAT computing center.
- This work has been carried out within the framework of the EUROfusion Consortium and has received funding from the Euratom research and training programme 2014-2018 under grant agreement No 633053. The views and opinions expressed herein do not necessarily reflect those of the European Commission.