



# ELECTRON THERMAL DIFFUSIVITY DIFFERENCES DUE TO MAGNETIC GEOMETRY ON THE HSX STELLARATOR

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## MOTIVATION

Experimentally investigate thermal transport on the Helically Symmetric eXperiment (HSX), a Quasi-Helically Symmetric (QHS) stellarator

- HSX has considerable magnetic geometry flexibility owing to a set of auxiliary coils surrounding its main field coils
- A Thomson Scattering system collects electron temperature and density profiles passively throughout experimental runs
- Changes in diffusive heat flux can be observed by the Thomson system when the magnetic geometry is altered
- An enhancement in heat flux has been measured when magnetic well depth is decreased by adding toroidal flux

## DIFFUSIVE HEAT FLUX MODEL

Electron temperature  $T_e$  and density  $n_e$  profiles are required to calculate an experimental electron thermal diffusivity  $\chi_e$  profile. The HSX Thomson system provides:

- $T_e, n_e$  at 10 spatial points over radius
- $n_e$  obtained from signal intensity
- $T_e$  obtained by Maximum Likelihood
- 6 individual discharges ensembled

Start with energy conservation for electrons:

$$\frac{\partial}{\partial t} \left[ \frac{3}{2} n_e T_e \right] + \nabla \cdot \mathbf{q}_e = p_{abs} - p_{loss} = p_{net}$$

Make the following assumptions:

- Plasma has reached a steady state
- Heat flux is diffusive  $\mathbf{q}_e = n_e \chi_e \nabla T_e$  within  $0.1 < r/a < 0.6$
- Absorbed power density is dominated by ECRH source
- Absorbed power density far exceeds losses  $p_{abs} \gg p_{loss}$

Energy conservation becomes:

$$\nabla \cdot \mathbf{q}_e = p_{abs} = p_0 \left[ 1 - (r/a)^2 \right]^{p_1}$$

where  $p_0$  is used to scale to absorbed power measured by diamagnetic loop at turn-off,  $p_1$  used to fit profile shape from raytracing.

Finally integrating over plasma volume yields:

$$q_e = \frac{1}{r} \int_0^r p_{abs} r' dr'$$

$$\chi_e = \frac{q_e}{n_e \nabla T_e}$$

## GAUSSIAN PROCESS REGRESSION (GPR) METHOD

If  $n_e$  and  $\nabla T_e$  profiles can be determined from Thomson data,  $\chi_e$  can be directly calculated. GPR provides:

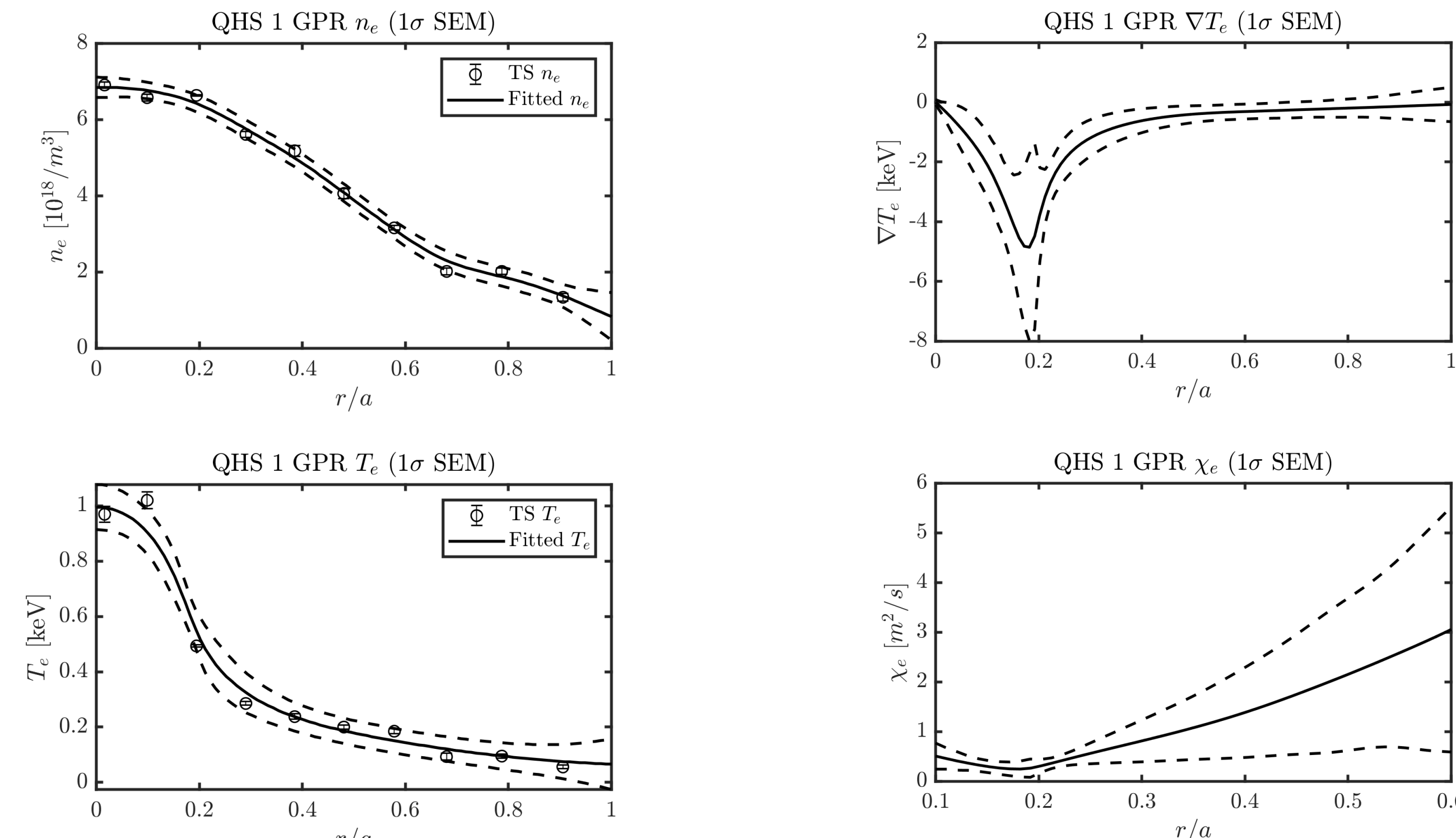
- Non-parametric smooth curve fitting
- Gradient fit constraints and predictions
- Bayesian-probability based method

GPTools library V0.2.3 [1] used for fitting,

- $n_e$  fit = SE kernel + linear mean + diagonal noise
- $T_e$  fit = SE kernel + hyperbolic tangent mean + diagonal noise

Problem - expanding uncertainty caused by singularity in error propagation as  $\nabla T_e \rightarrow 0$ :

$$\sigma_{\chi_e} = \chi_e \sqrt{\left( \frac{\sigma_{n_e}}{n_e} \right)^2 + \left( \frac{\sigma_{\nabla T_e}}{\nabla T_e} \right)^2 + \left( \frac{\sigma_{q_e}}{q_e} \right)^2}$$



## LEAST SQUARES (LSQ) REGRESSION METHOD

To avoid calculating  $\nabla T_e$  directly, instead use an assumed form for  $\chi_e$  [2]:

$$\chi_e = \chi_0 + \chi_1 r + \chi_2 r^2$$

then calculate a modeled  $T_{mdl}$  which can be fit to  $T_e$  measurements by nonlinear Least Squares:

$$T_{mdl} = \int_0^r \frac{\partial T_e}{\partial r'} dr' + T_{peak}(r=0)$$

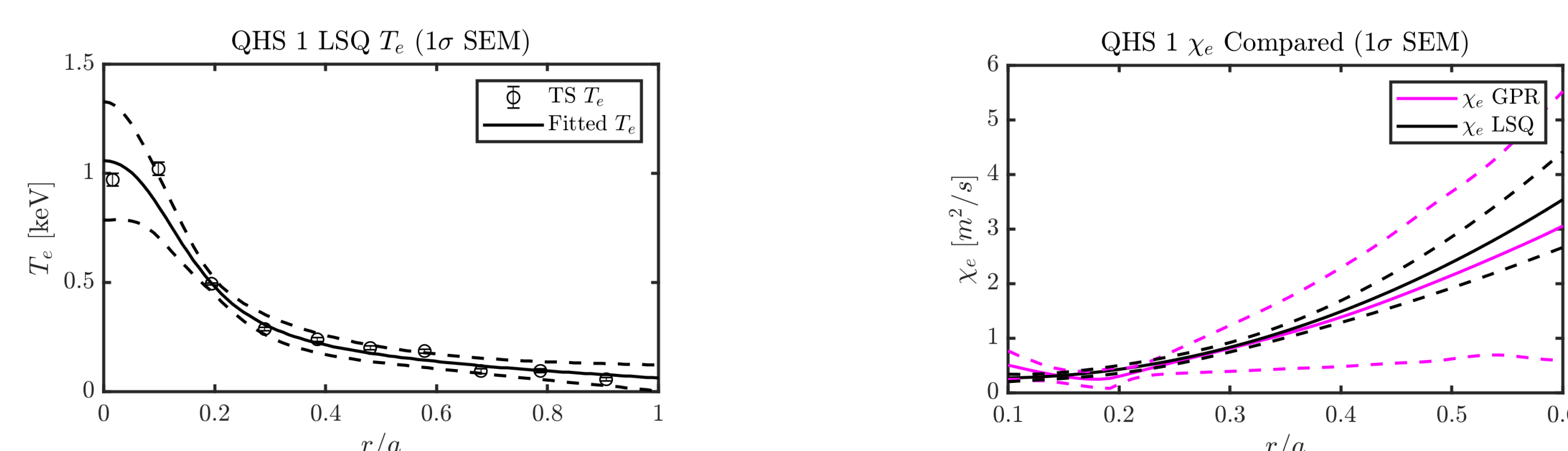
$$F_i = \left( T_e(r_i) - T_{mdl}(r_i) \right) \frac{1}{\sigma_{T_e}(\rho_i)}$$

$$\mathbf{X} = [T_{peak}, \chi_0, \chi_1, \chi_2]$$

Uncertainty in  $\chi_e$  is then determined based upon minimization parameters via covariance matrix, calculated using weighted Jacobian and Mean Squared Error (MSE):

$$\mathbf{Cov} = \text{inv}(\mathbf{J}' * \mathbf{J}) * \text{MSE}$$

$$\sigma_{\chi_e}^2 = \chi_e^2 \left( \left( \frac{\sigma_{n_e}}{n_e} \right)^2 + \left( \frac{\sigma_{q_e}}{q_e} \right)^2 \right) + \mathbf{Cov} * \mathbf{r}^T \mathbf{r}$$



## HSX MAGNETIC CONFIGURATIONS

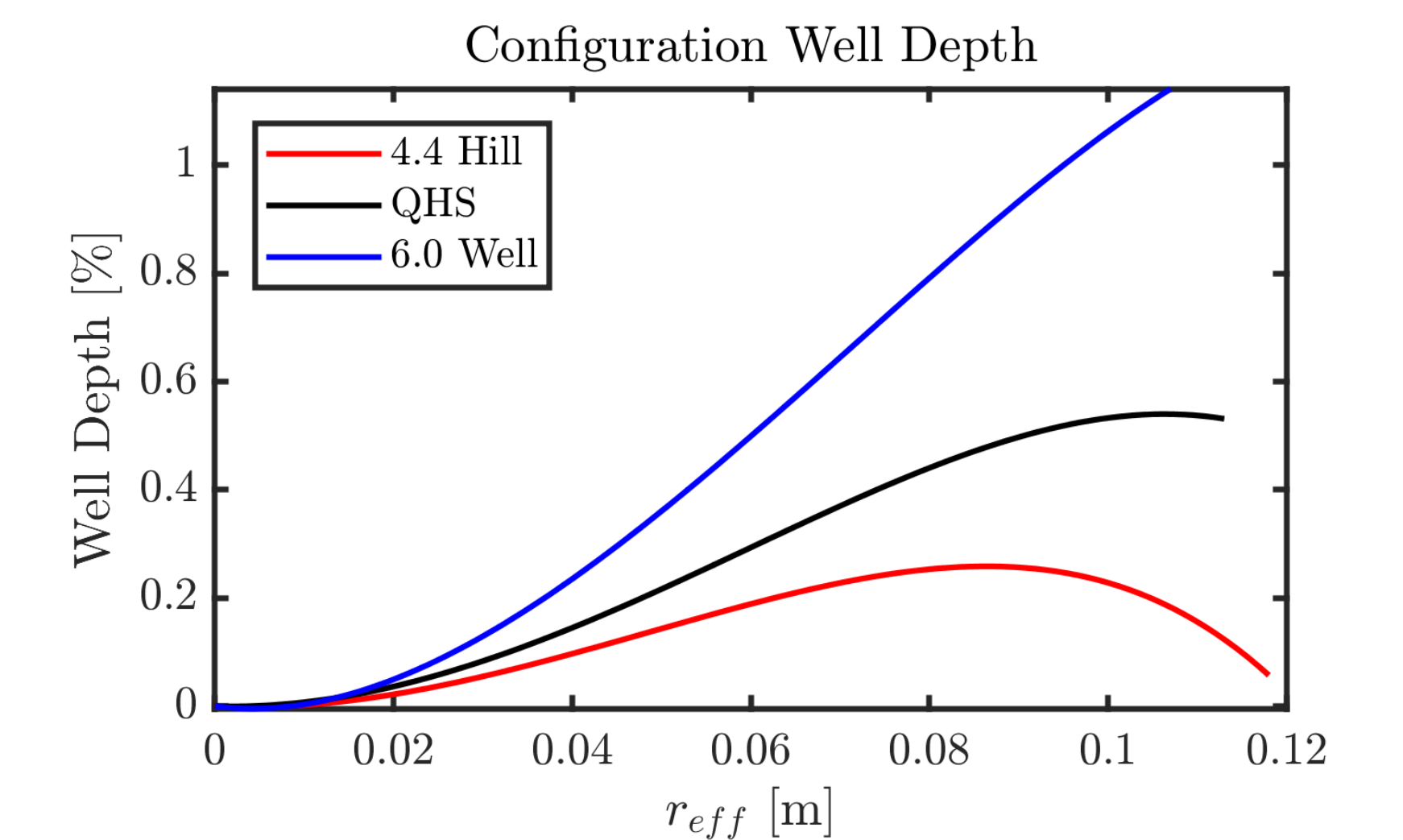
For toroidal flux  $\Psi$  and volume  $V$ , relative well depth  $W_D$  is:

$$W_D = \frac{dV/d\Psi(r/a=0) - dV/d\Psi(r/a)}{dV/d\Psi(r/a=0)}$$

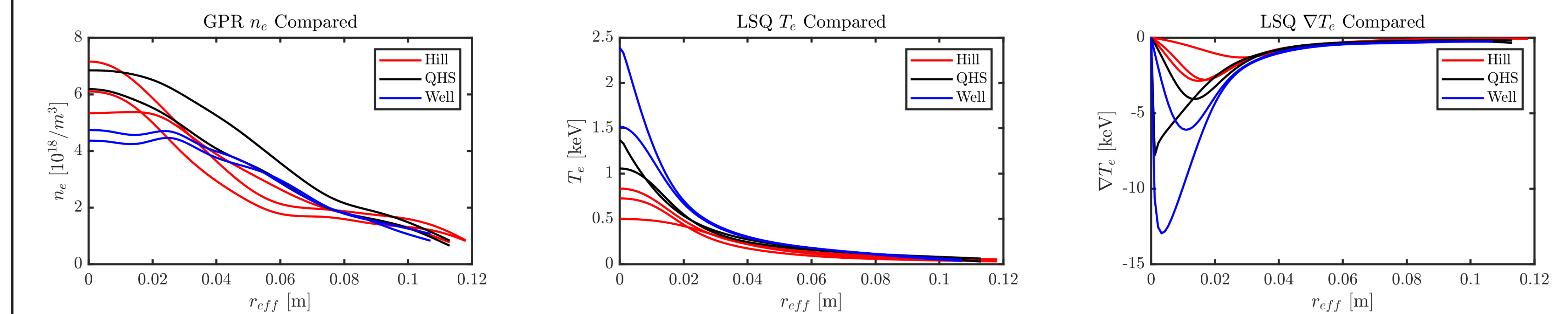
Hill - QHS - Well Configuration range of HSX geometry:

- HILL - increased  $\Psi$ , decreased  $W_D$ , increased plasma radius
- QHS - neoclassically optimized geometry
- WELL - decreased  $\Psi$ , increased  $W_D$ , decreased plasma radius

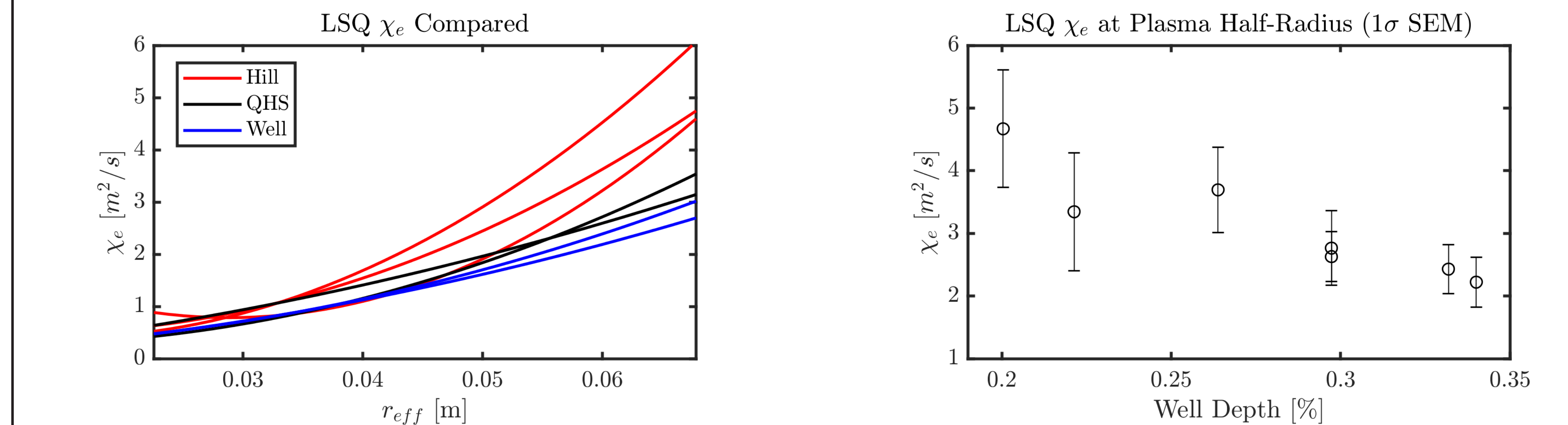
Configurations labeled by auxiliary coil current fraction.



## FITS USED FOR FINAL RESULTS



## DIFFUSIVITY RESULTS



EXPERIMENTAL  $\chi_e$  TREND: HILL (SMALL  $W_D$ ) > QHS > WELL (LARGE  $W_D$ )

## SUMMARY & FUTURE WORK

Summary:

- A TS system has been used to make electron density and temperature measurements over a range of magnetic geometries
- These  $n_e$  and  $T_e$  profiles have been used to calculate electron thermal diffusivity using a diffusive energy flux model
- An enhanced diffusivity with decreasing magnetic well may suggest increased turbulent transport in those configurations

Future Work:

- Modeling will be performed to attempt to separate turbulence from neoclassical transport
- The mechanism primarily responsible for changing transport with respect to magnetic well will be identified
- The TS system will be upgraded to improve signal digitization and analysis techniques

## REFERENCES

- [1] M. A. Chilenski, M. Greenwald, Y. Marzouk, N. T. Howard, A. E. White, J. E. Rice, and J. R. Walk. Improved profile fitting and quantification of uncertainty in experimental measurements of impurity transport coefficients using Gaussian process regression. *Nuclear Fusion*, 55(2):23012, 2015.
- [2] H. Maassberg, R. Burhenn, U. Gasparino, G. Kühner, H. Ringler, and K. S. Dyabilin. Experimental and neoclassical electron heat transport in the LMFP regime for the stellarators W7-A, L-2, and W7-AS. *Physics of Fluids B*, 1993.

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