



# Collisionless zonal flow damping in quasisymmetric stellarators





J. Smoniewski<sup>1</sup>, E. Sánchez<sup>2</sup>, I. Calvo<sup>2</sup>, M.J. Pueschel<sup>3</sup>, J.N. Talmadge<sup>1</sup>

<sup>2</sup>Laboratorio Nacional de Fusión, CIEMAT, Madrid, Spain

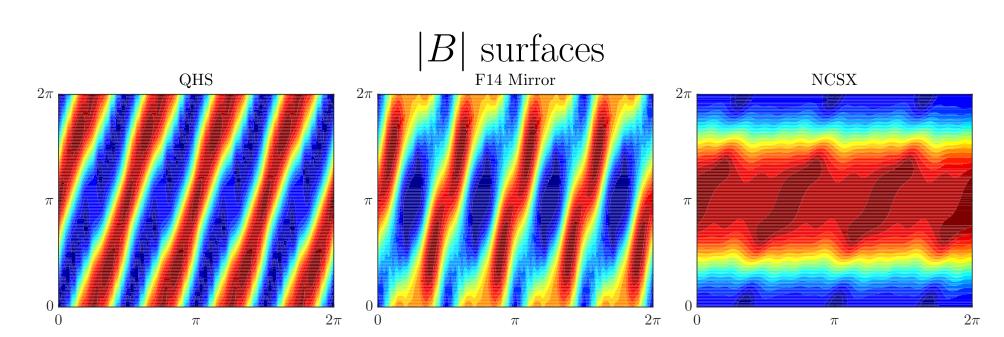
<sup>1</sup>HSX Plasma Laboratory, University of Wisconsin-Madison

#### MOTIVATION

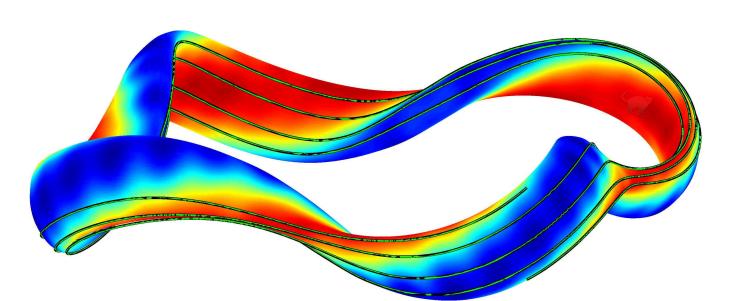
- Optimization for turbulent transport requires predictive capabilities.
- Zonal flows are often important to turbulence saturation, and thus transport.
- Collisionless zonal flow damping simple linear calculation for an optimization loop.
- Can a linear zonal flow (ZF) damping calculation predict the turbulence heat flux?
- -Not the residual alone, but perhaps damping/oscillation.
- ZF residual finite as  $k_x \rho_s \to 0$  in tokamaks, vanishes in non-axisymmetry (Monreal 2016).
- Are real quasi-symmetric stellarators similar to tokamaks (finite  $R_{\rm ZF}$  as  $k_x \rho_s \to 0$ )?
- -No,  $R_{\rm ZF} \to 0$  as  $k_x \rho_s \to 0$  if any radial particle drift.

#### I. HSX and NCSX configurations

- $\bullet$  **HSX** auxiliary coils  $\rightarrow$  allows variation of radial drift (neoclassical transport), important for ZF oscillation.
- QHS Quasi-Helical Symmetry, configuration optimized for reduced neoclassical transport and flow damping.
- F14 Mirror Broken symmetry with [n,m] = 4.0 and 8.0mirror term, effective ripple like a conventional stellarator.



• NCSX – Quasi-axisymmetry, three period device with a dominant [n,m] = 0,1 mode. Large bootstrap current similar to a tokamak.



- Gyrokinetics: GENE (www.genecode.org) & EUTERPE -Gene local flux tube, with 1-8 poloidal turns.
- -Gene full flux surface, real space poloidal discretization. -Euterpe fully global domain.
- No unstable modes, no gradient drive, no collisions, adiabatic electrons.
- Initialize a  $k_y = 0$ , finite- $k_x$  zonal mode and calculate time evolution.

#### II. Zonal flow damping in a stellarator

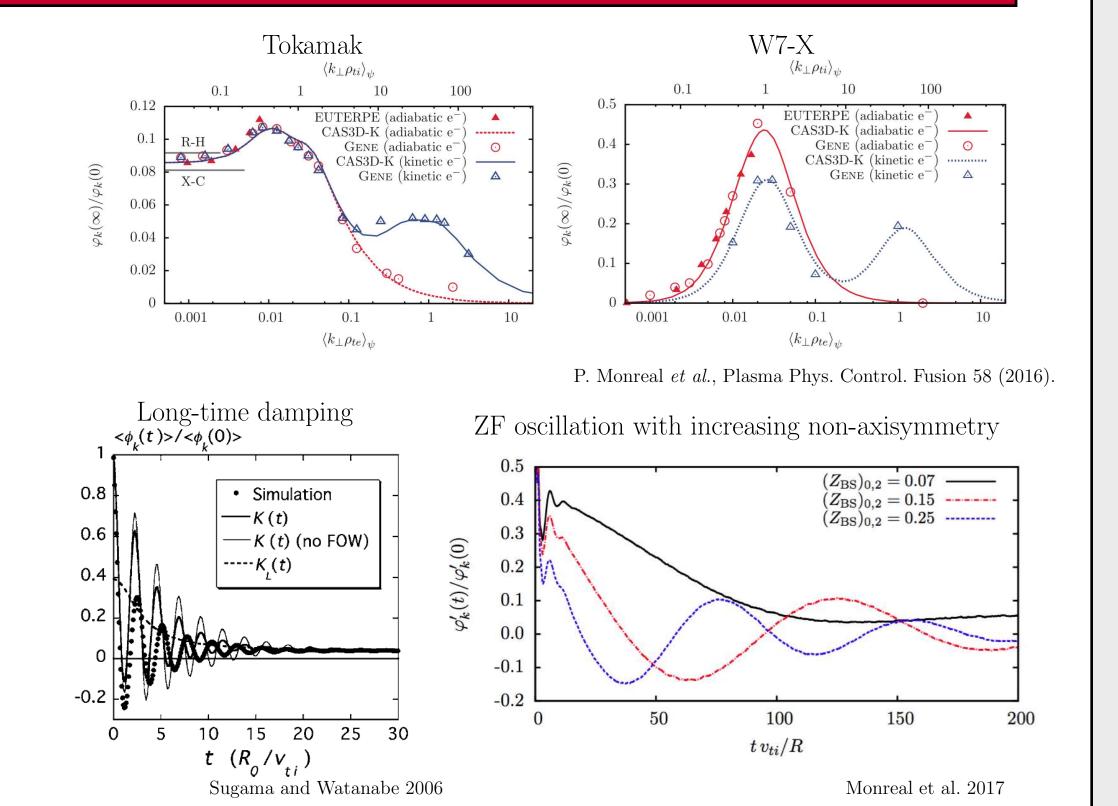
- Zonal Flow (ZF) poloidally symmetric potential perturbation, drives flows by  $E \times B$  drift.
- Flows shear turbulent eddies, reduce radial correlation length, and regulate transport.
- Given initial impulse, long time residual may relate to ability to support zonal flows.
- In a circular tokamak, Rosenbluth-Hinton zonal flow residual  $R_{\rm ZF}$  finite as  $k_x \rho_s \to 0$ .:

$$\lim_{t \to \infty} \frac{\phi(t)}{\phi(0)} = \frac{1}{1 + 1.6q^2/\epsilon^{1/2}}$$

• Tokamak with q = 1.5 and  $\epsilon \approx \frac{a}{R} \approx 0.13$ ,

$$\lim_{t \to \infty} \frac{\phi(t)}{\phi(0)} = 0.09$$

- In Wendelstein 7-X,  $R_{\rm ZF} \to 0$  at small  $k_x \rho_s$ .
- Differences for non-axisymmetry: long-time damping and zonal flow oscillations.
- GAMs quickly damped for low-shear stellarators, and not discussed here.



- Long-time damping related to radial drift,  $\tau_c \sim 1/|k_r \bar{v}_{dr}|$ .
- ZF oscillation present if Landau damping small enough, and frequency  $\Omega_{\rm ZF}$  increases with radial drift.
- Optimized devices minimize neoclassical loss, reduce long-time damping, reduce frequency of zonal flow oscillations.

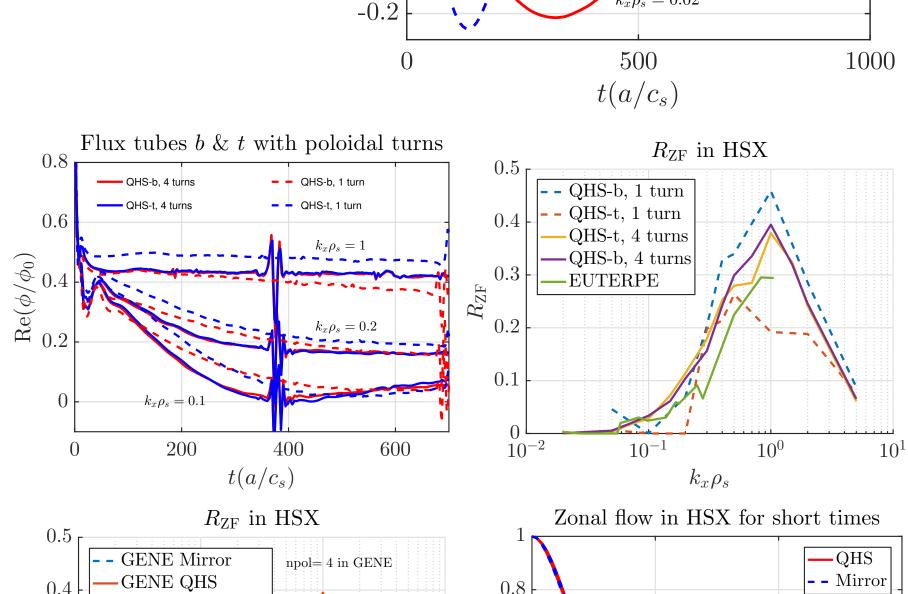
Zonal flow in HSX for long times

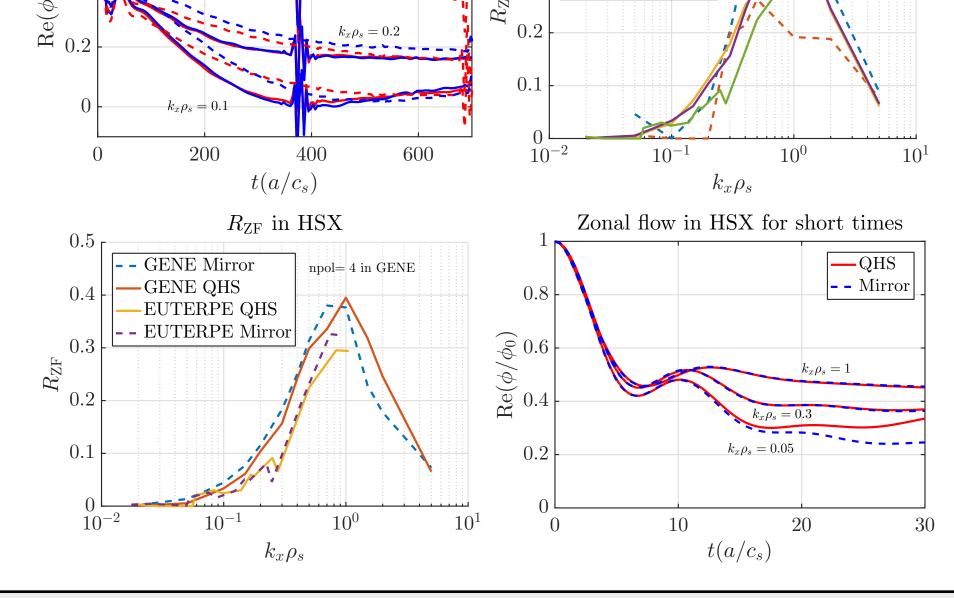
### III. Realistic HSX configurations match other stellarators

- Do non-ideal quasi-symmetric devices have finite  $R_{\rm ZF}$  as  $k_x \rho_s \to 0$ ?
- $\lim_{t \to \infty} \frac{\phi(t)}{\phi(0)} = 0.64$ • HSX-equivalent tokamak,  $q \to q_{\text{eff}} = \frac{1}{|m\iota - n|} = 1/3$ ,  $\epsilon \to \epsilon_h \approx 0.14 (r/a) \approx 0.1$
- Real HSX,  $R_{\rm ZF} \to 0$  as  $k_x \rho_s \to 0$ ,  $R_{\rm ZF}$  is smaller than Rosenbluth-Hinton estimate.
- ZF oscillation parameters found by nonlinear least squares fit.
- Fit ZF oscillation amplitude  $A_{\rm ZF}$ , frequency  $\Omega_{\rm ZF}$ , and damping  $\gamma_{\rm ZF}$ , ZF residual  $R_{\rm ZF}$ , and algebraic damping.

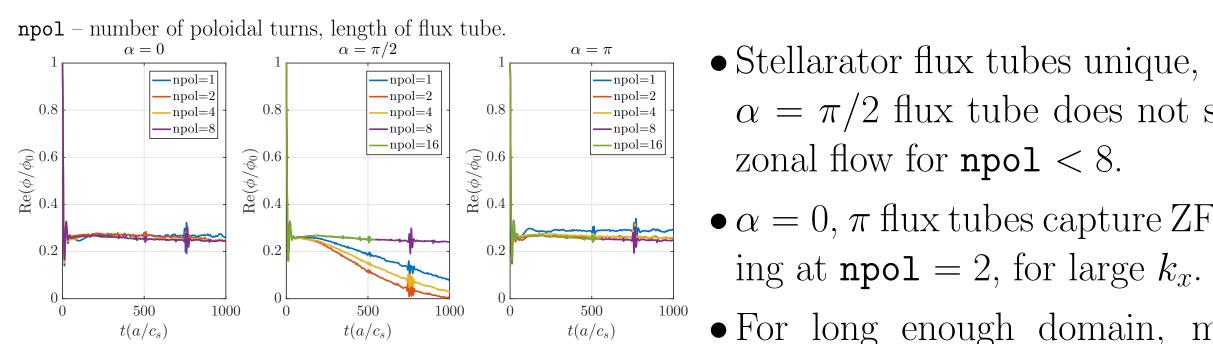
$$\varphi_k'(t)/\varphi_k'(0) = A_{\rm ZF}\cos(\Omega_{\rm ZF}t)\exp(-\gamma_{\rm ZF}t) + R_{\rm ZF} + \frac{c}{1+dt^e}$$

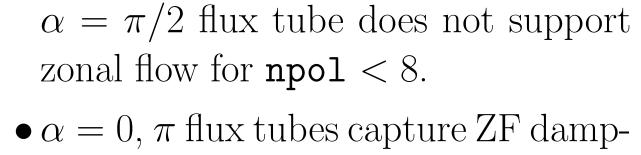
- Zonal flow oscillations present at low- $k_x \rho_s$ . Frequency decreases with reduced radial drifts from neoclassical optimization.
- Two stellarator symmetric flux tubes for local Gene calculations. QHS-b(ean) and QHS-t(riangle) centered in good and bad curvature, respectively.
- Damping depends on flux surface averages  $\rightarrow$  long flux tube (> 1 poloidal turn) required to match different flux tubes and global EUTERPE.
- ullet No significant difference in  $R_{\mathrm{ZF}}$ QHS/Mirror. Large difference in  $\Omega_{\rm ZF}$ .
- For short times comparable to turbulence correlation time – no difference in QHS/Mirror.
- Heat flux larger in F14 Mir- 🗟 ror, despite matched  $R_{\rm ZF}$ . (Nonlinear simulation in flux tube with 4 poloidal turns)

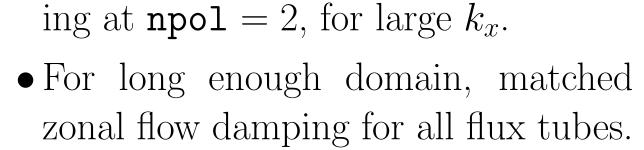


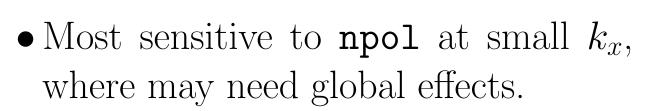


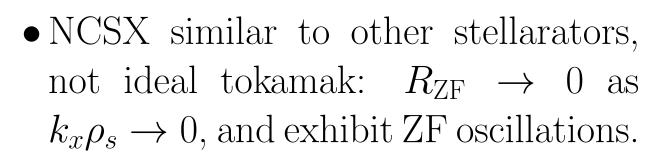
#### IV. Zonal flow residual in NCSX

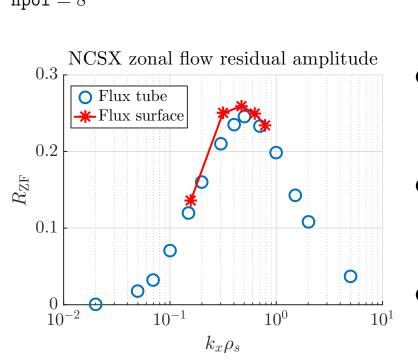






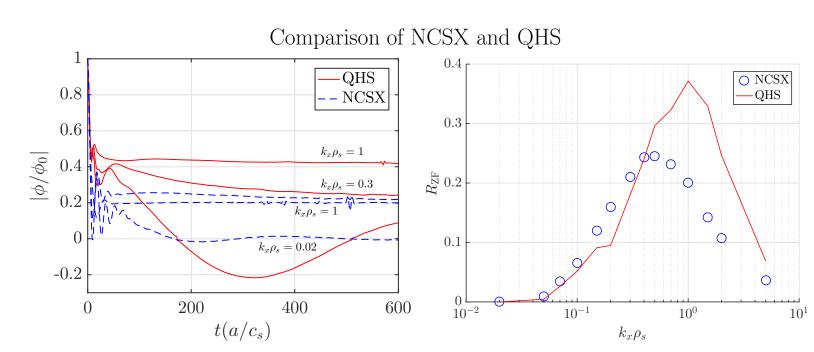






<sup>3</sup>University of Texas at Austin

- Full flux surface  $R_{\rm ZF}$  in agreement with flux tube calculations for npol = 8. Global simulations in progress.
- Agreement of flux tube and full surface suggests communication between flux tubes not important for residual.
- Zonal flow oscillations only present for very small  $k_x$ , amplitude much smaller than QHS.



ullet Peak  $R_{\mathrm{ZF}}$  and peak  $k_x$ smaller than QHS. Requires further consideration.

## V. Summary

- Zonal flow damping requires flux surface averages sensitive to geometry representation.  $R_{\rm ZF}$  captured by local calculation only with extended flux tube.
- Poloidal turn requirements depend on geometry: some flux tubes do not support zonal flows unless sampling enough of the surface.
- Realistic, as opposed to ideal, quasi-symmetric devices differ from tokamaks with  $R_{\rm ZF} \to 0$  as  $k_x \rho_s \to 0$  and finite ZF oscillation frequency.
- QHS, compared to Mirror, exhibits reduced oscillation frequency and decay consistent with reduced radial particle drifts in optimized stellarators.
- Trends between configurations that are not captured by linear instability dynamics cannot be explained by  $R_{\rm ZF}$  behavior alone.

## Acknowledgments

- This work is supported by US DOE Grant DE-FG02-93ER54222, and used resources of the National Energy Research Scientific Computing Center, a DOE Office of Science User Facility supported by the Office of Science of the U.S. Department of Energy under Contract No. DE-AC02-05CH11231.
- This work has been partially funded by the Ministerio de Economía y Competitividad of Spain under project ENE2015-70142-P. The authors thankfully acknowledge the computer resources, technical expertise and assistance provided by the Barcelona Supercomputing Center-Centro Nacional de Supercomputación and the CIEMAT computing center.
- This work has been carried out within the framework of the EUROfusion Consortium and has received funding from the Euratom research and training programme 2014-2018 under grant agreement No 633053. The views and opinions expressed herein do not necessarily reflect those of the European Commission.