

Overview of Recent Results from HSX

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HSX Plasma Laboratory

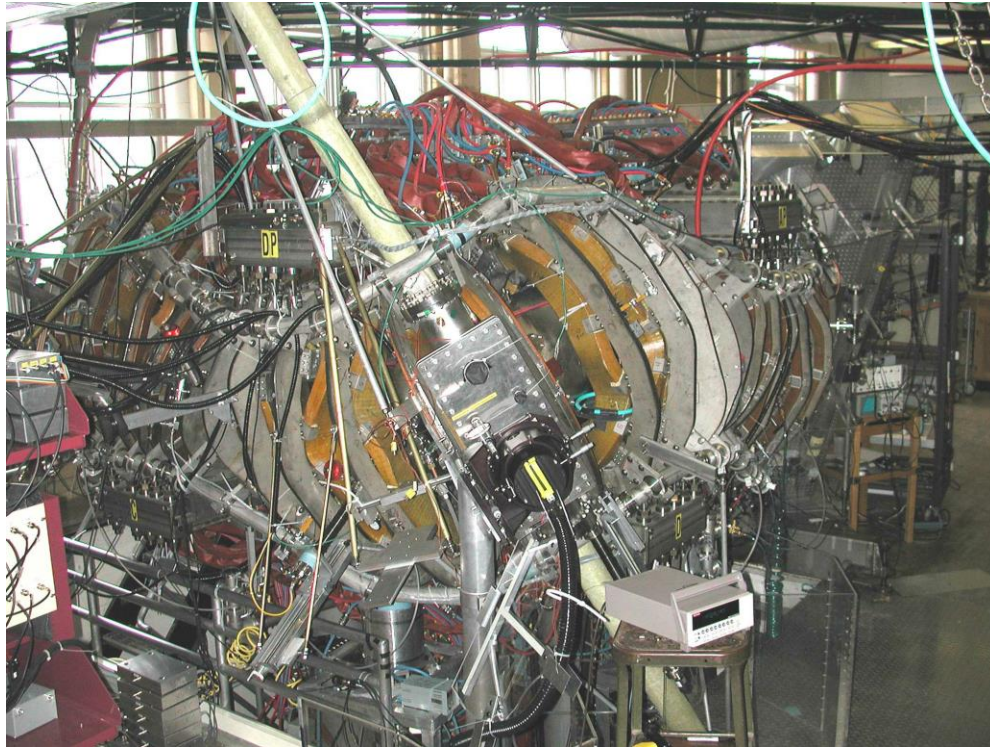
The University of Wisconsin-Madison

14th International Stellarator Workshop

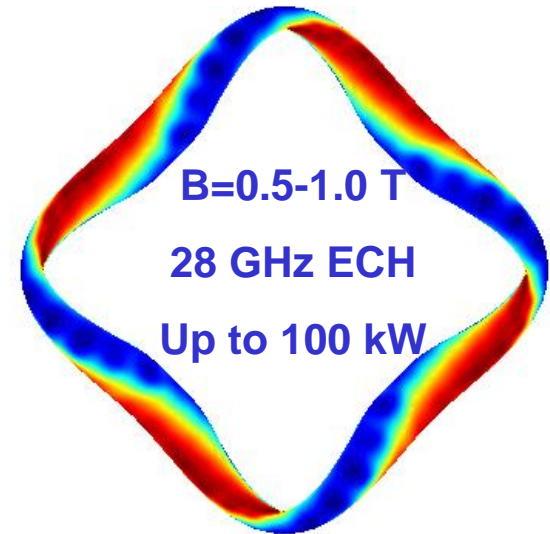
Greifswald, Germany

Sept. 22-26, 2003

HSX is a Quasi-helically Symmetric Stellarator

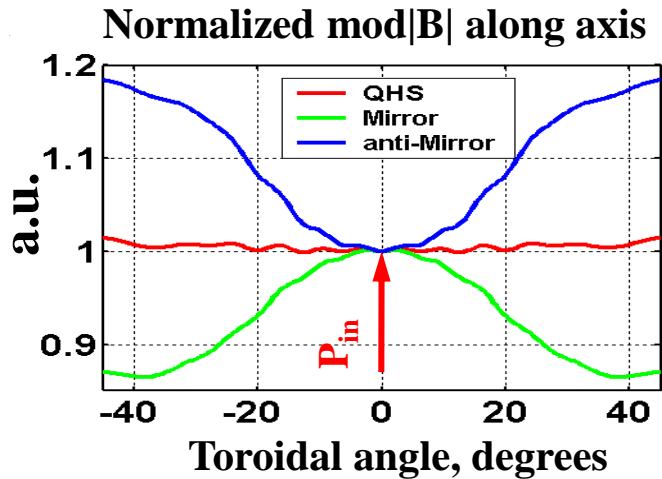


Magnitude of B



HSX has a helical axis of symmetry in $|B|$ and a resulting very low level of neoclassical transport

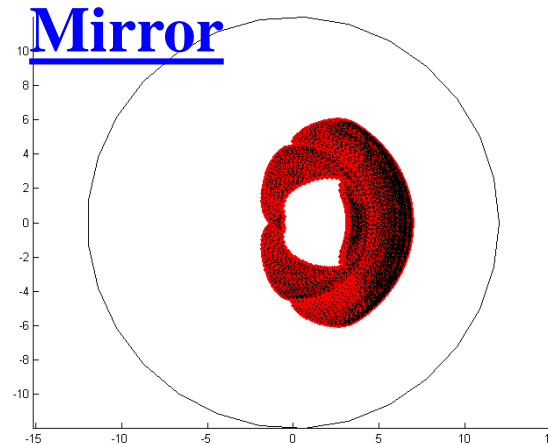
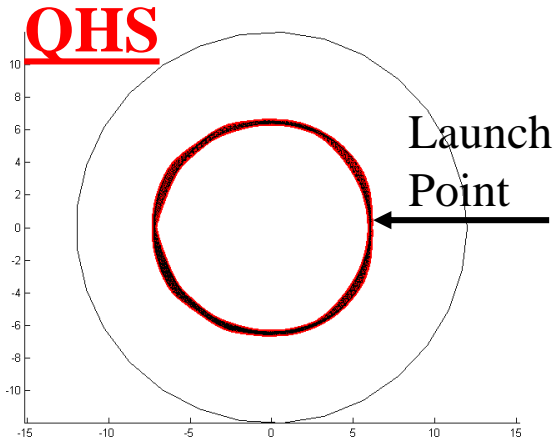
Neoclassical Transport Can Be Increased with Mirror Field



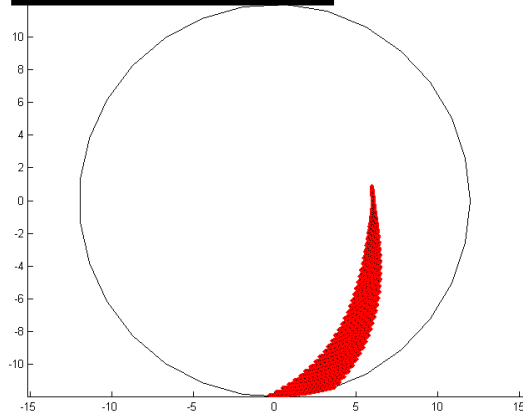
- Mirror configurations in HSX are produced with auxiliary coils in which an additional toroidal mirror term is added to the magnetic field spectrum

- In Mirror mode the term is added to the main field at the location of launching antenna
- In anti-Mirror it is opposite to the main field

Trapped Particle Orbits



anti-Mirror



- Trapped particles in QHS are well-confined
- By the ECH antenna, orbits are poor in Mirror configuration;
- Even worse in anti-Mirror

Anomalous Transport Should Dominate Thermal Plasmas Under Present Operation

- How does anomalous transport scale in HSX?
 - Evidence that X_e scales like $1/n$ (vs $T_e^{3/2}$)
- In lower density operation strong evidence for energetic tail population
 - Well-confined in QHS
- What are the benefits of QHS in more thermal plasmas?
 - Good absorption of ECH
 - Reduced rotation damping
 - Eventually, good confinement of thermal plasmas in $lmfp$ regime

ASTRA is Used to Model Transport

- In addition to the neoclassical transport, we also assume an anomalous electron thermal conductivity:

$$\chi_e = \chi_{e,neo} + \chi_{e,anom}$$

- Previously, we assumed an anomalous thermal conductivity based on ASDEX L-mode scaling:

$$\chi_{e,anom} \sim \frac{T_e^{3/2}}{RB^2} \frac{1}{\left[1.1 - (r/a)^2\right]^4}$$

- If $\tau \sim 1/T_e^{3/2} = nT/P$, then:

$$T \sim (P/n)^{0.4}; \quad \tau \sim (n/P)^{0.6}; \quad W \sim n^{0.6}P^{0.4}; \quad \text{ISS95-like}$$

Alcator-like Model Fits Data Better

- ASDEX L-mode model did not agree with scaling dependencies of experimental data.
- A better model of anomalous transport in HSX is an Alcator-like dependency (n_e in units of 10^{18} m^{-3}):

$$\chi_{e,anom} = \frac{10.35}{n_e} \text{ m}^2 / \text{s}$$

- If $\tau \sim n = nT/P$, then:

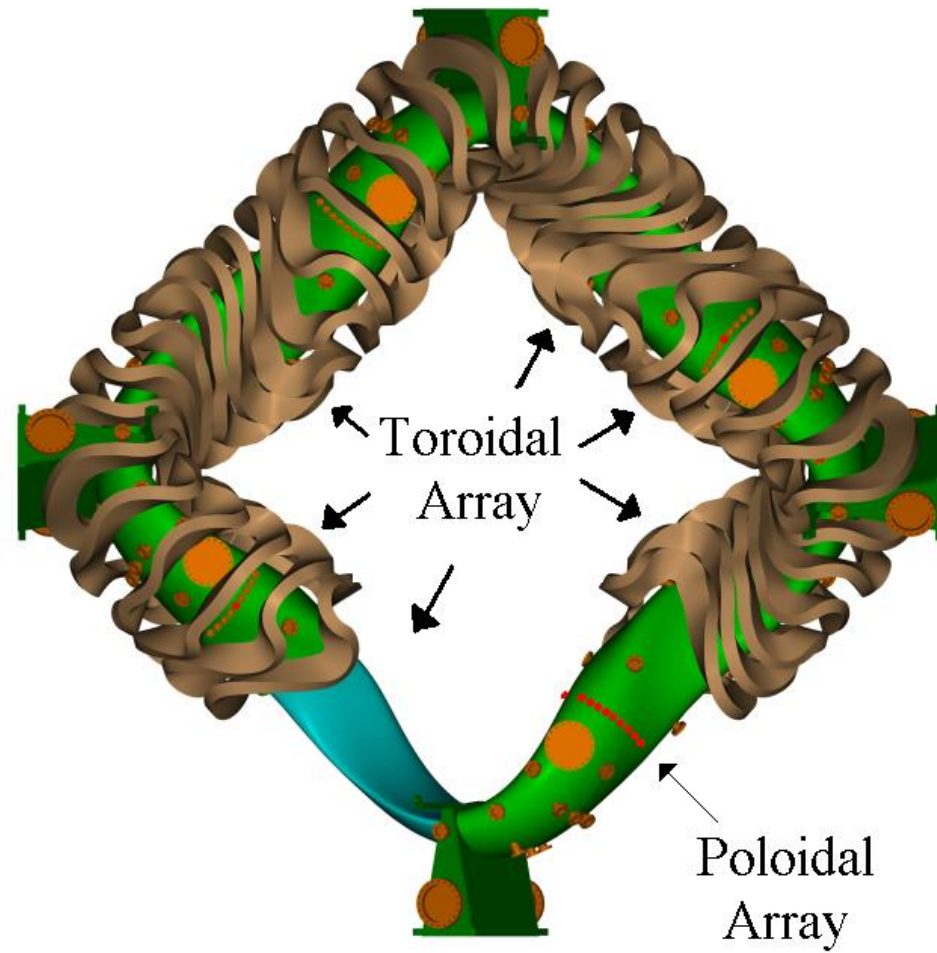
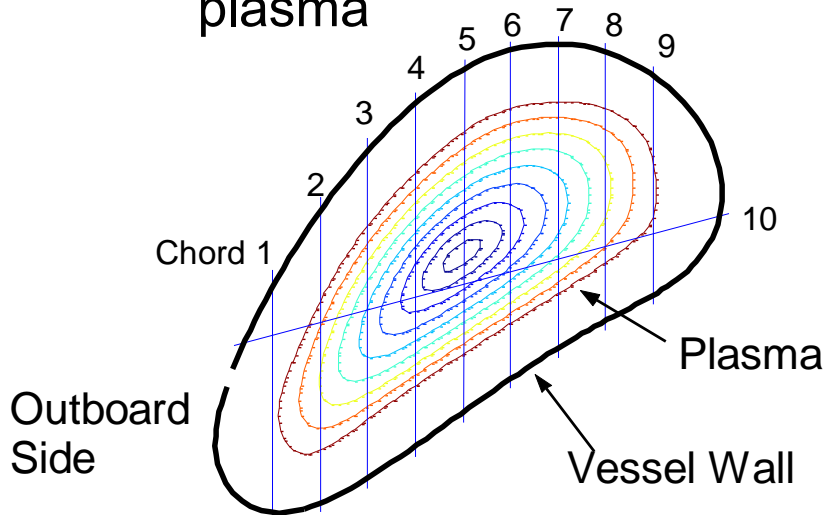
$T \sim P$ (independent of n) ; $\tau \sim n$; $W \sim nP$;

which is more in agreement with experiment

Full details on transport modeling and experimental measurements this afternoon in poster by Talmadge

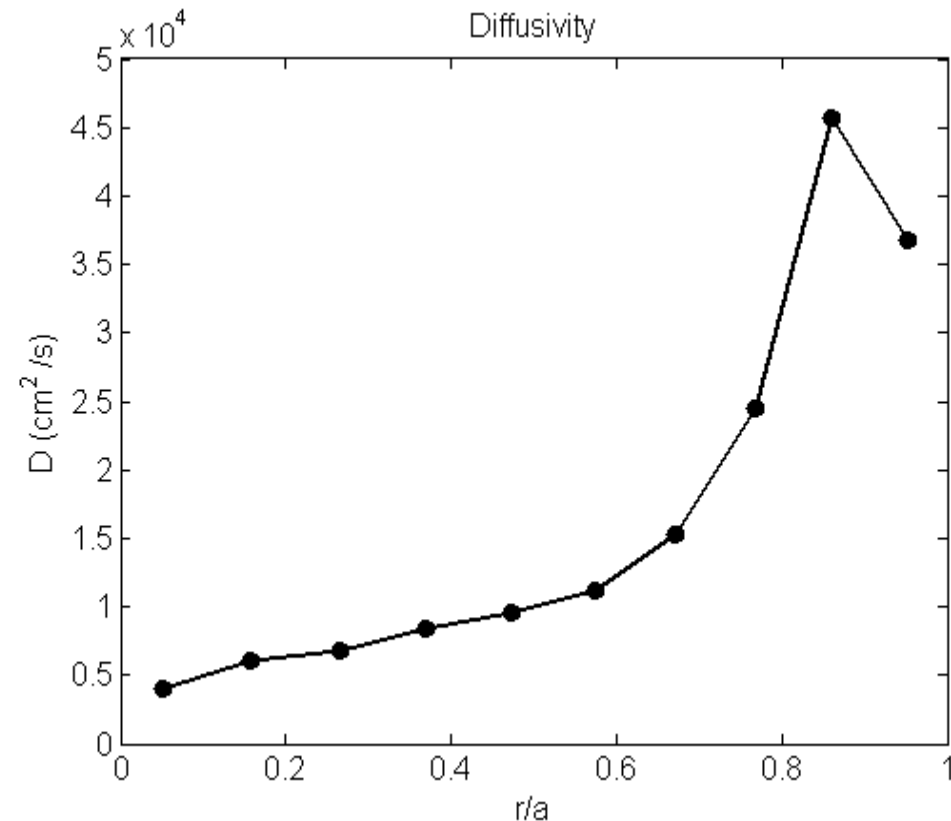
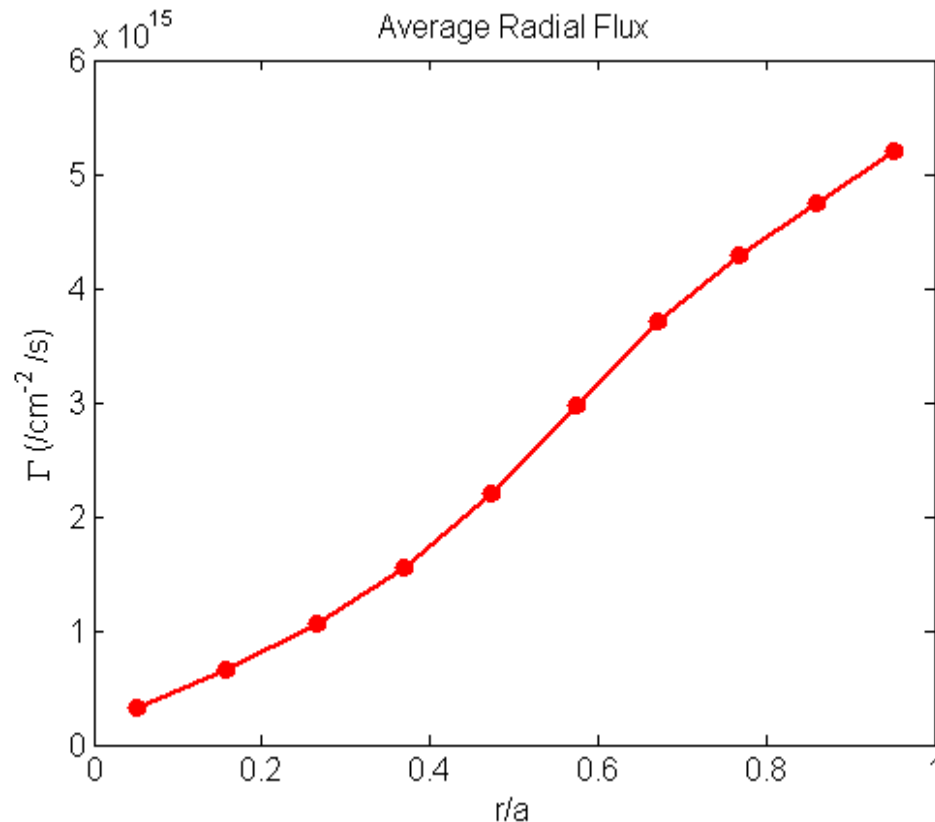
H_α Array Used to Measure Neutral Density

- HSX has 16 H_α detectors forming two arrays
 - Toroidal array: 7 detectors on magnetically equivalent ports
 - Poloidal array: 9 detectors viewing cross section of plasma

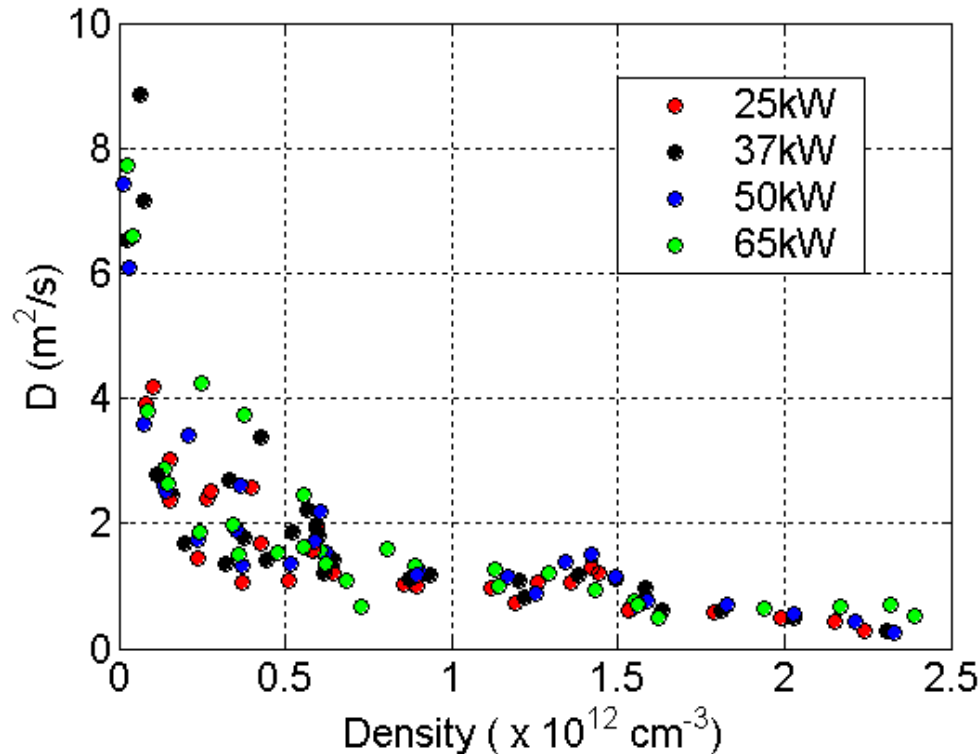


Estimate of Diffusion Coefficient Give $\sim 1 \text{ m}^2/\text{s}$

- The total source inside each flux surface and the electron density gradient give an effective diffusion coefficient
- Yields $D(r) \sim 10^4 \text{ cm}^2/\text{s}$, increasing towards edge
- $n_e = 1.5 \times 10^{12} \text{ cm}^{-3}$, $P = 37 \text{ kW}$



H_α Measurements Consistent with Model

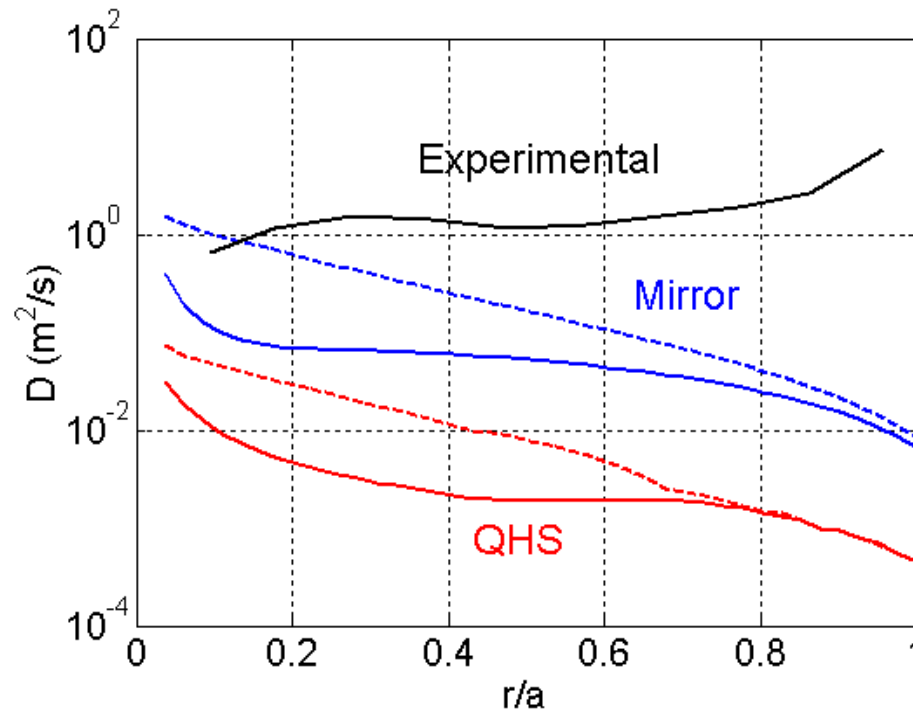


- Negligible dependence on power!

- See poster by J. Canik this afternoon
- H_α toroidal and poloidal data analyzed using DEGAS code for 3 different line average densities and 4 different power levels
- Dependence of diffusion coefficient on n and P :

$$D_{anom} \sim \frac{P^{0.09}}{n^{0.6}}$$

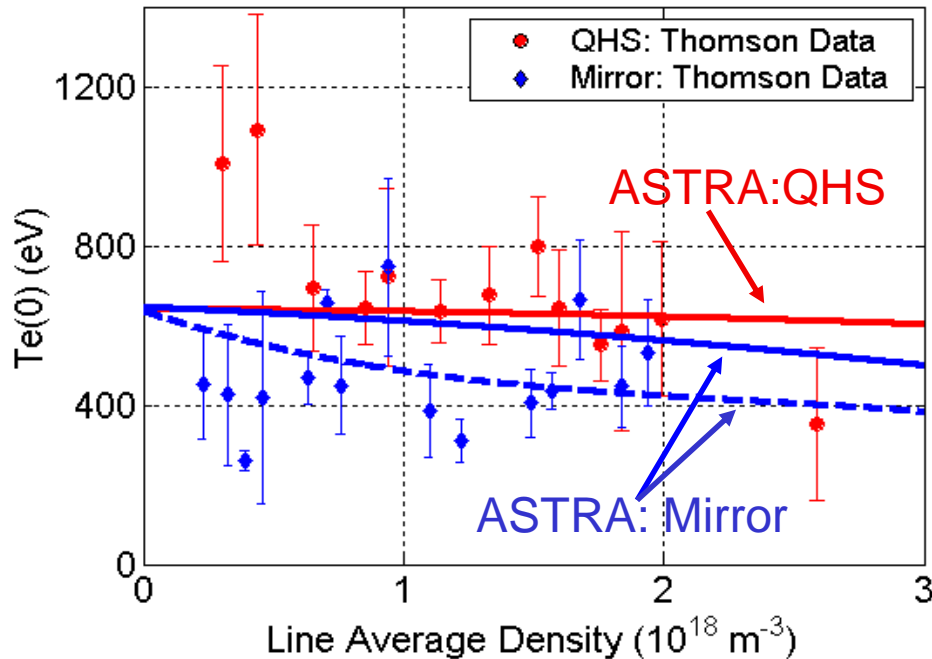
Experimental Diffusion Coefficients Larger than Neoclassical Values



ASTRA calculations of neoclassical diffusion coefficients with ambipolar E_r (solid) and $E_r = 0$ (dashed)

D from equilibrium analysis in rough agreement with modulated gas puff

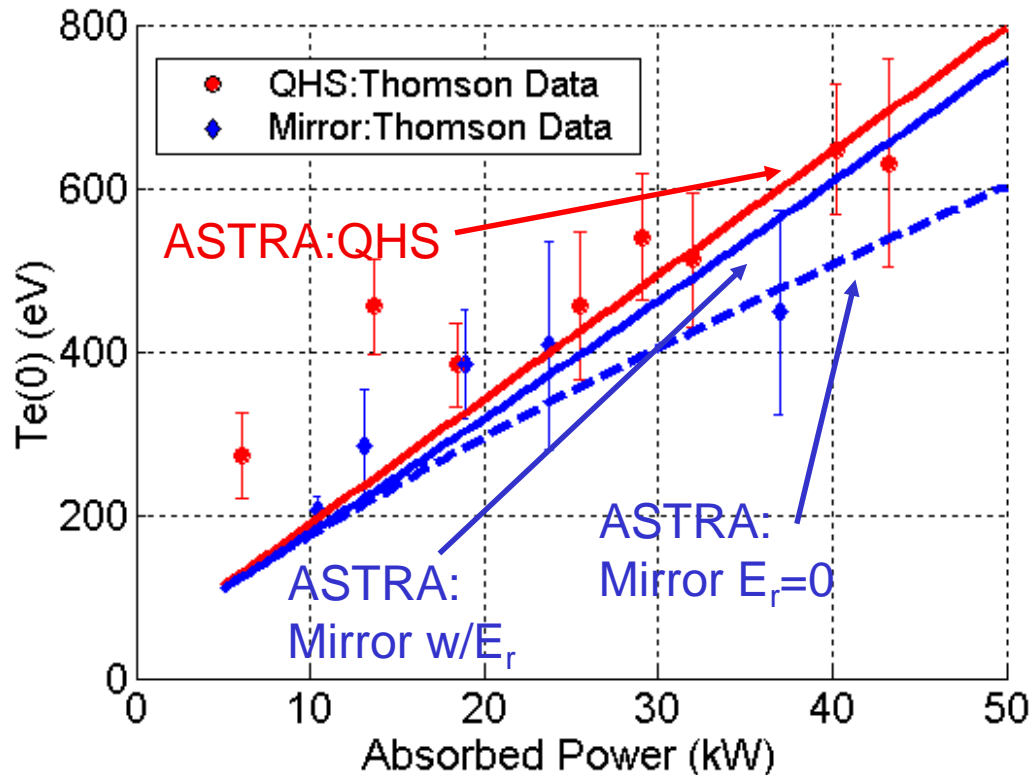
Central Electron Temperature is Independent of Density



- QHS thermal conductivity is dominated only by anomalous transport
- $T_e(0)$ in Mirror is calculated with self-consistent E_r (solid line) and $E_r = 0$ (dashed).

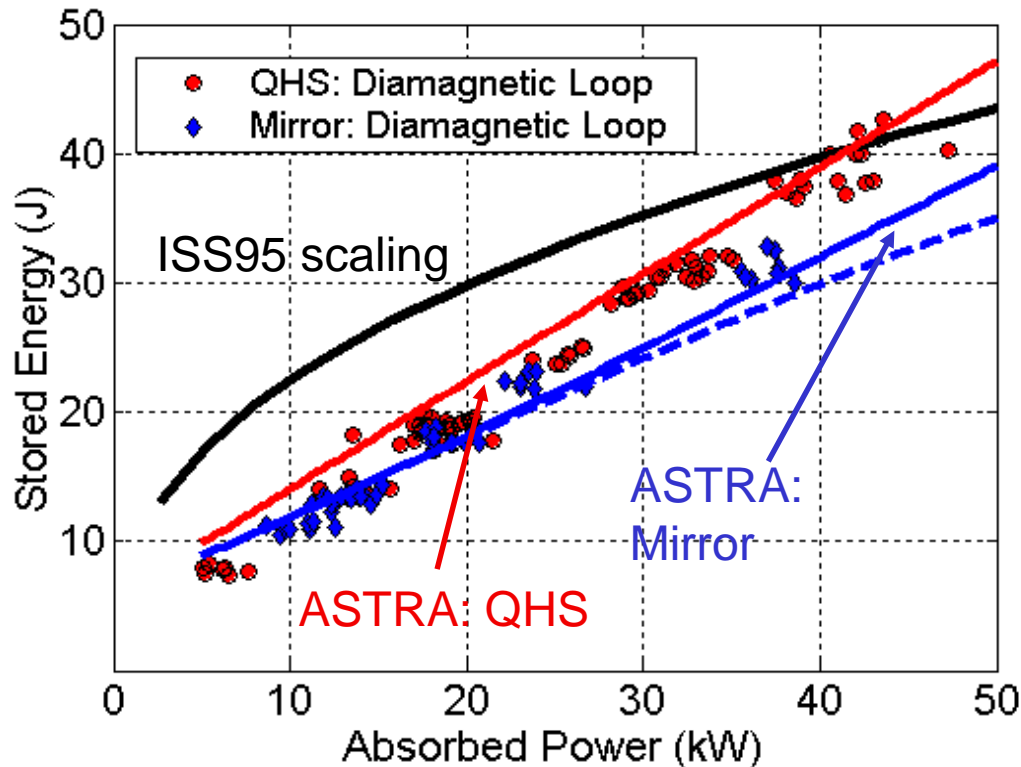
- Except for lowest densities, $T_e(0)$ from Thomson scattering is roughly independent of density,
- Consistent with $\chi \sim 1/n$ model.

Thomson Data shows $T_e(0)$ Increases Linearly with Power



- Fixed density of $1.5 \times 10^{18} \text{ m}^{-3}$.
- ASTRA calculation is consistent with Thomson measurements for QHS and Mirror
- $T \sim P$ is supportive of $\chi \sim 1/n$ model.

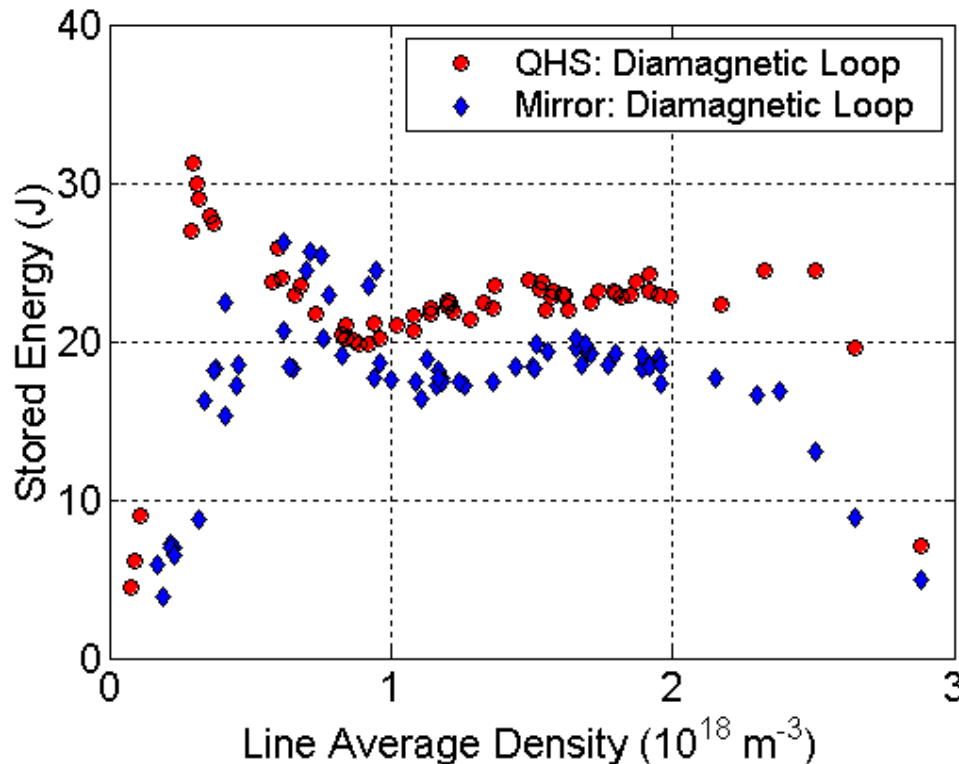
Stored Energy Increases Linearly with Power



- Fixed density of $1.5 \times 10^{18} \text{ m}^{-3}$.
- Difference in stored energy between QHS and Mirror reflects 15% difference in volume.
- $W \sim P$ in agreement with $\chi \sim 1/n$ model.

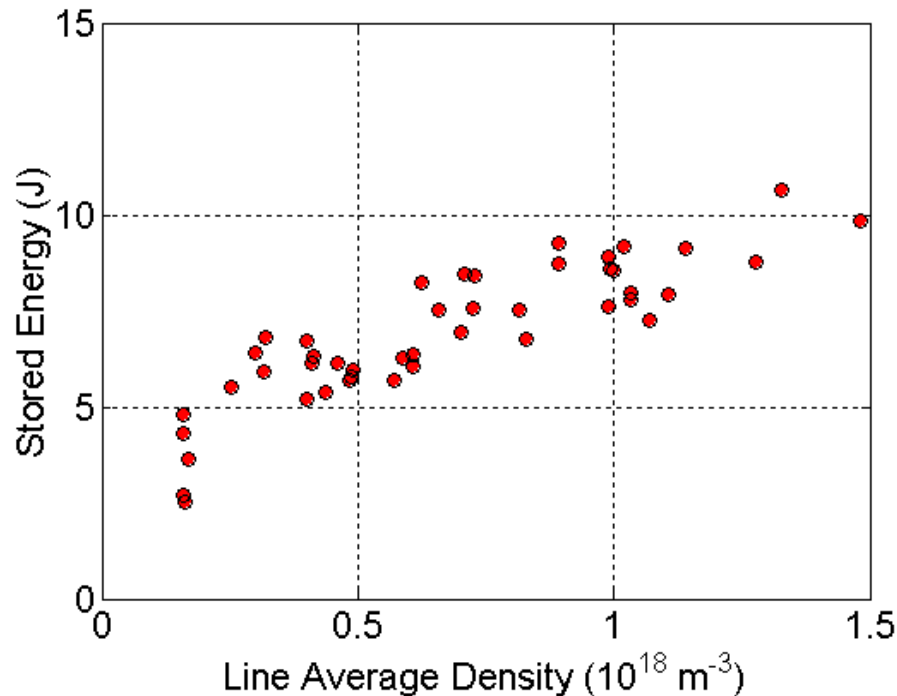
At the “Cross-roads”; higher-power and higher-density operation will be interesting!

Stored Energy Does Not Have Linear Dependence on Density



- Fixed input power, 40 kW.
- For $\chi \sim 1/n$ model, $W \sim n$ for fixed power. Data clearly does not show this.
- Are nonthermal electrons causing stored energy to peak quickly at low density?

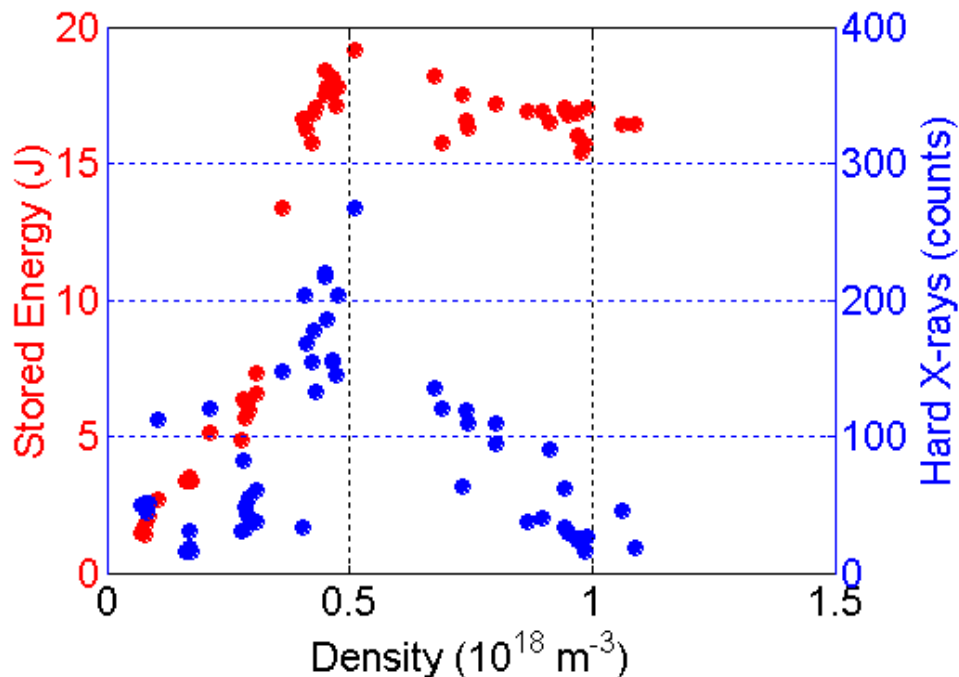
Stored Energy Goes Up Linearly with Density when Confinement is Poor



- Resonance is on low-field side of Mirror configuration where confinement of trapped particles is degraded
- $W \sim n$ in this configuration is now consistent with $\chi \sim 1/n$ model.

• Stored energy of 7 J at $n = 0.7 \times 10^{18} \text{ m}^{-3}$ now in agreement with ASTRA prediction

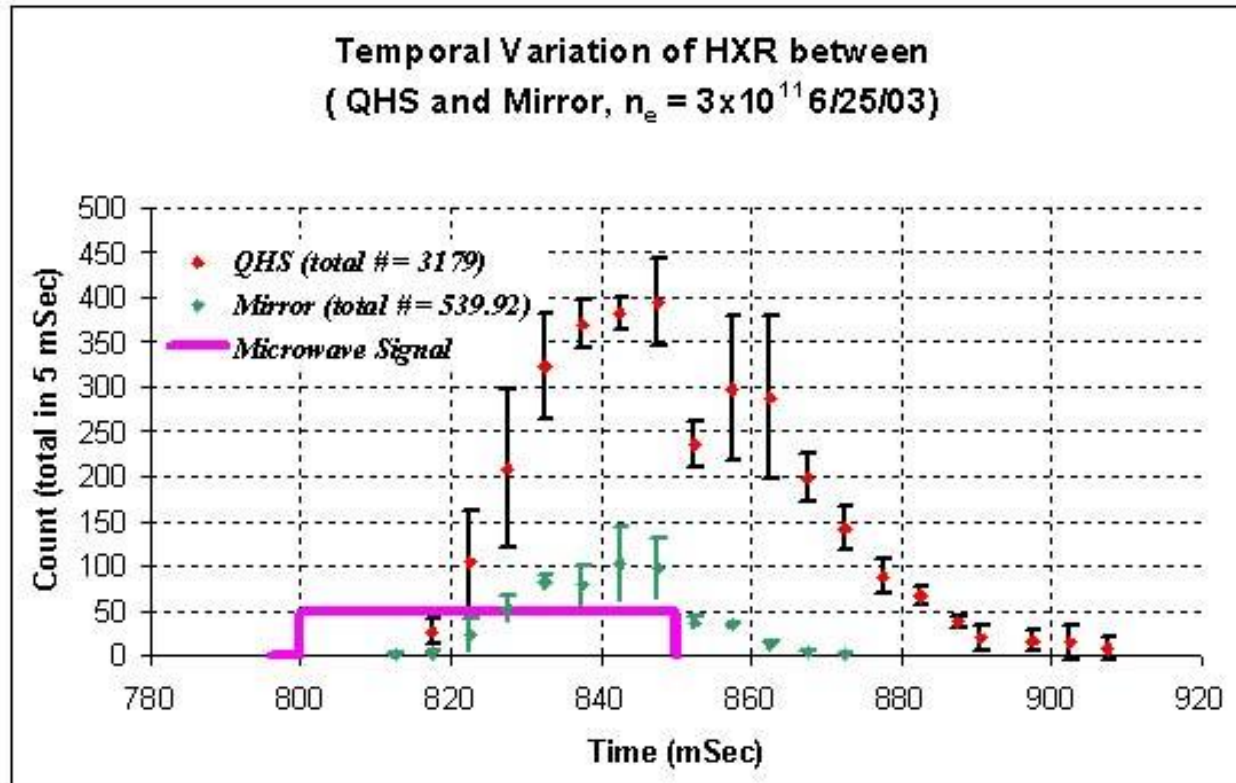
Hard X-rays Have Similar Dependence on Density as Stored Energy for $n < 0.5 \times 10^{18} \text{ m}^{-3}$



- Shielded and collimated CdZnTe detector with 200 μm stainless steel filter.
- Fixed input power: 40 kW.
- Hard X-ray intensity peaks at $0.5 \times 10^{18} \text{ m}^{-3}$, as does stored energy.

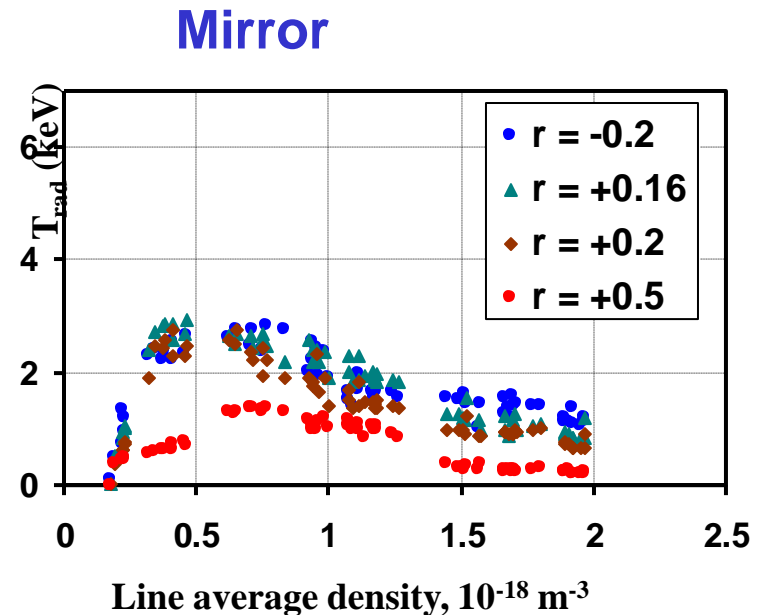
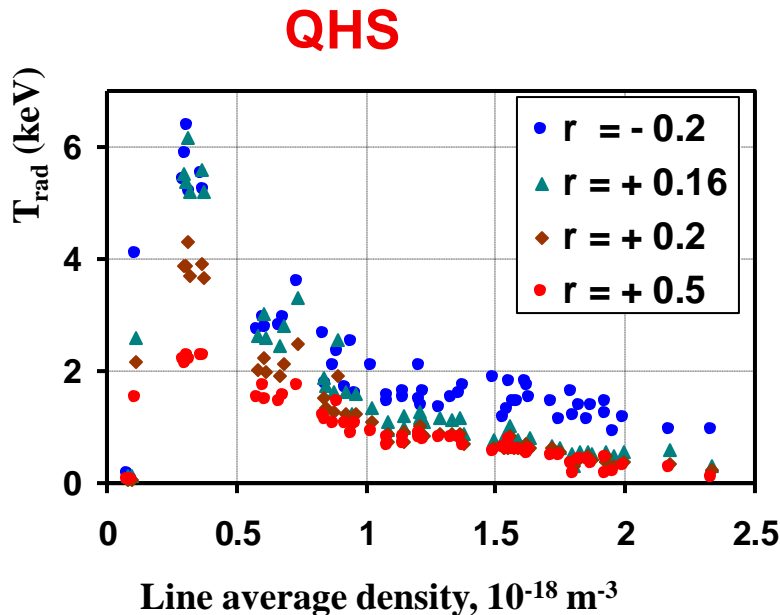
- Hard X-ray intensity falls off sharply beyond $1 \times 10^{18} \text{ m}^{-3}$, while stored energy remains roughly constant.

Hard X-rays Greater in QHS than Mirror



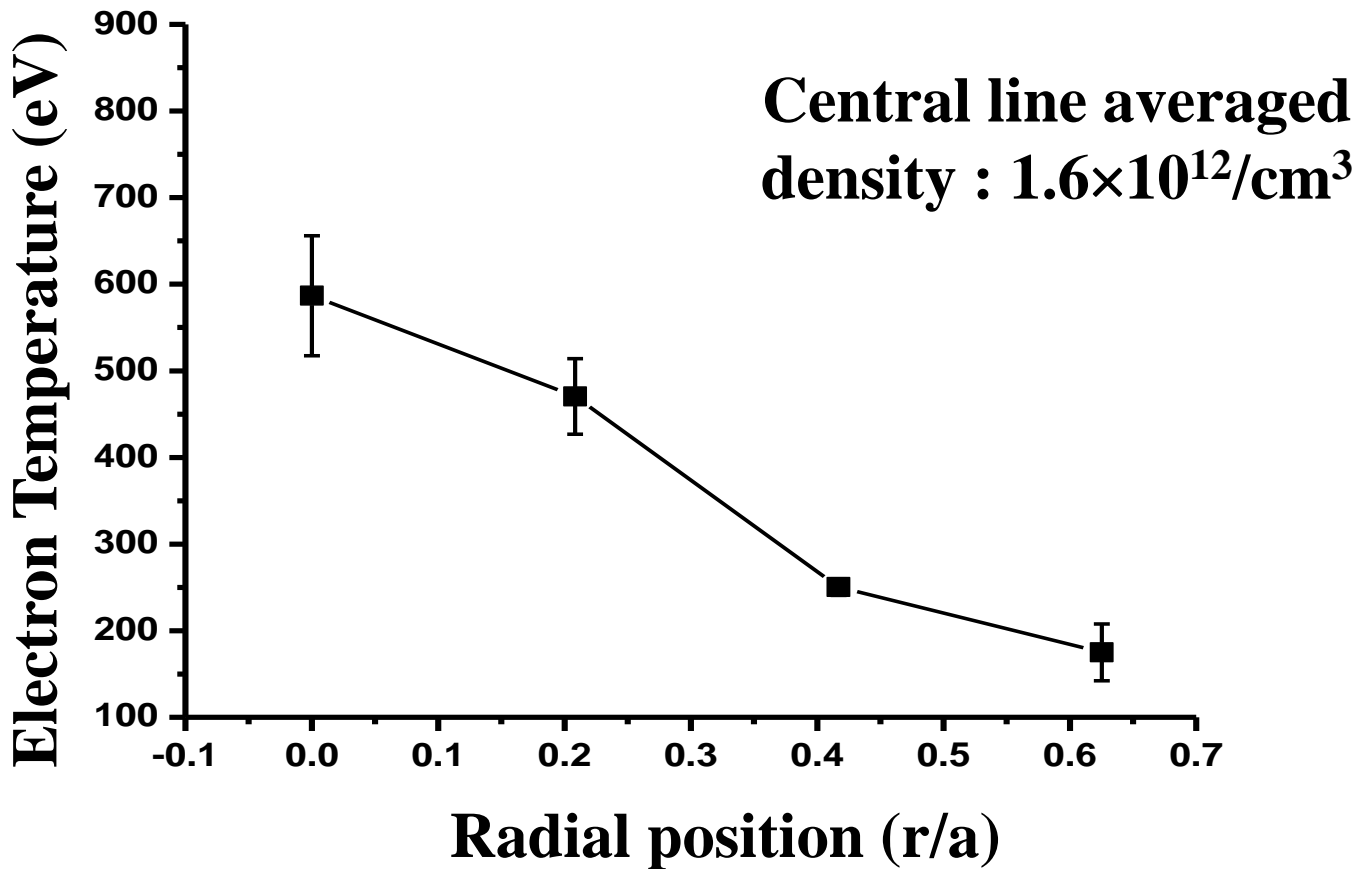
- Intensity increases till gyrotron turn-off, then decreases with 13 ms time constant for QHS, 5 ms for Mirror; virtually no hard x-ray counts for anti-Mirror

T_{rad} by ECE Shows Large Non-Thermal Component at Low Densities



- T_{rad} as high as 6 keV in **QHS**; less than 3 keV in **Mirror**
- As density increases emission approaches a thermal level in both QHS and Mirror plasmas
- Strong support for energetic tail contribution at low densities

QHS Central Resonance Thomson Scattering Profile

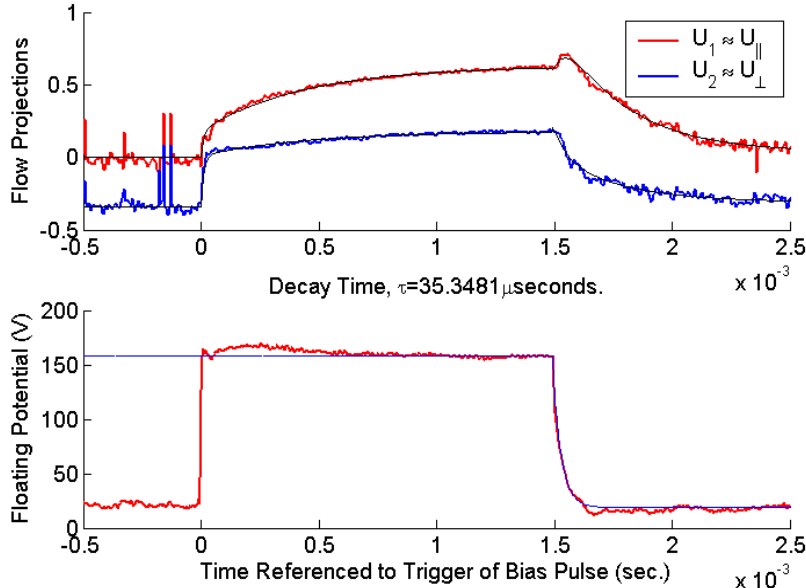


Thomson scattering shows reasonably peaked profiles
at higher densities with $T_{e0} \sim 600$ eV

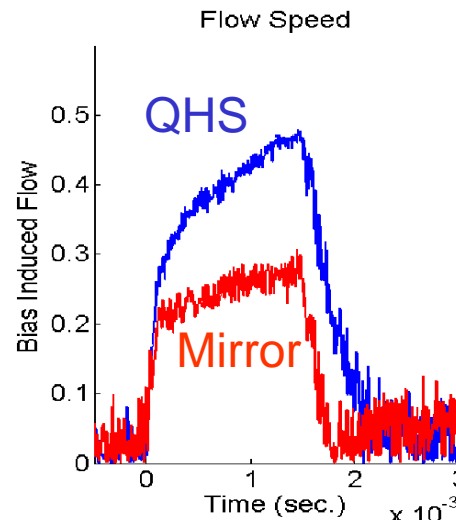
Differences Appear Between QHS and Mirror Modes for Thermal Plasmas

Biased electrode used to spin plasma

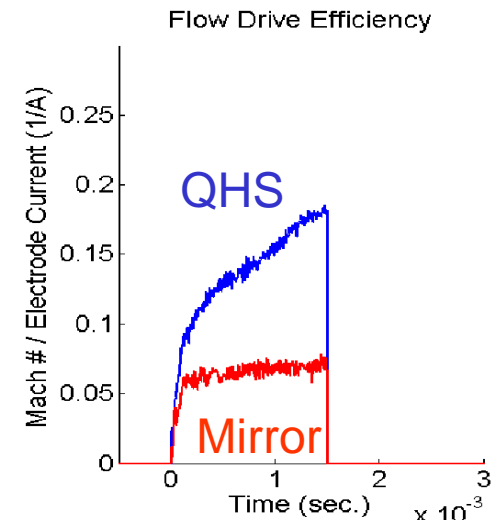
Fast Rise: 11.7139, Slow Rise: 438.3352, Fast Fall: 50.9474, Slow Fall: 278.6465 μ sec.



Two time scales for flow rise
(This is QHS case)

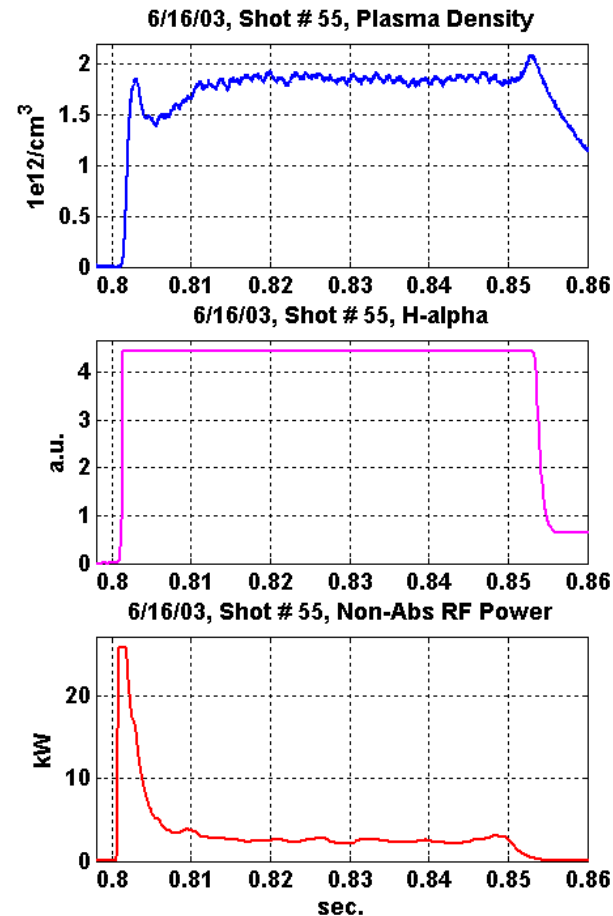
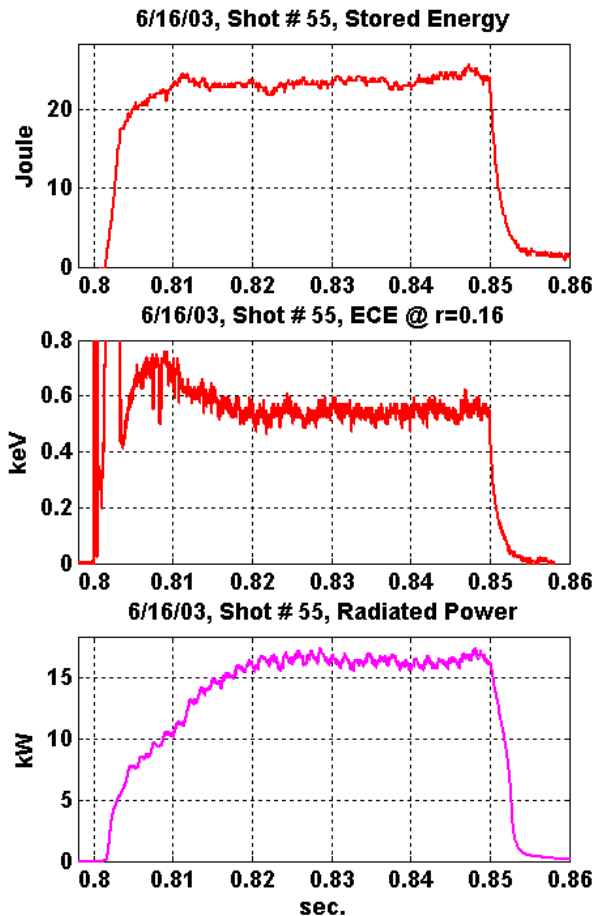


QHS flow damps slower, goes faster,
for less drive



Talk by Gerhardt this session

Stable, Thermal Discharges Are Achieved with QHS at Higher Density

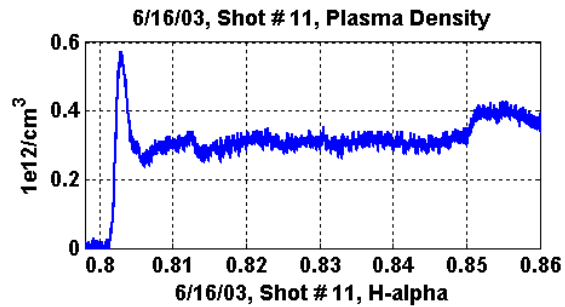
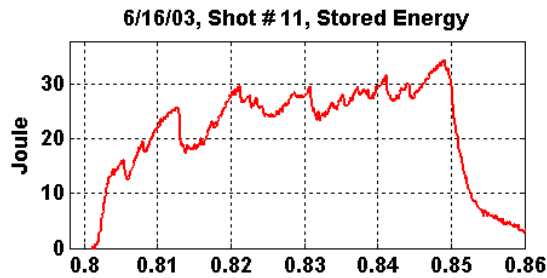
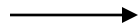


$T_e(0) \sim 600$ eV, line averaged density $\sim 1.8 \times 10^{18} \text{ m}^{-3}$

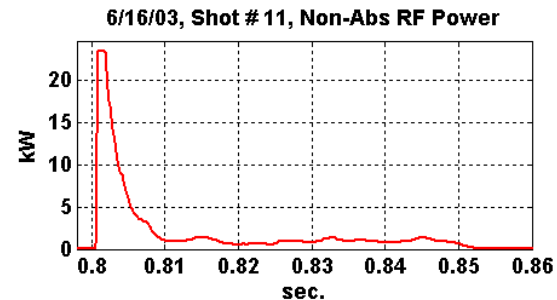
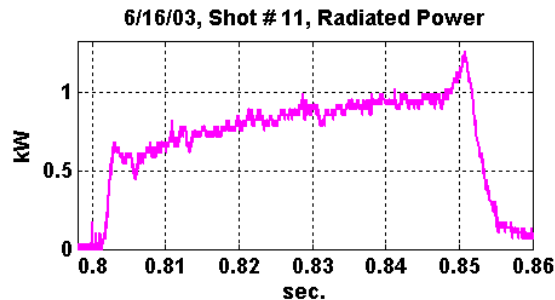
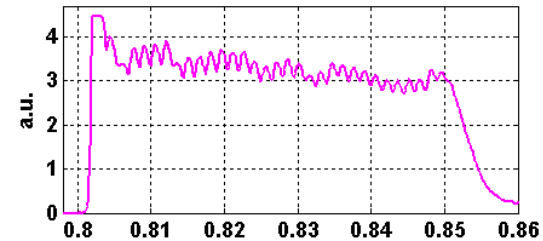
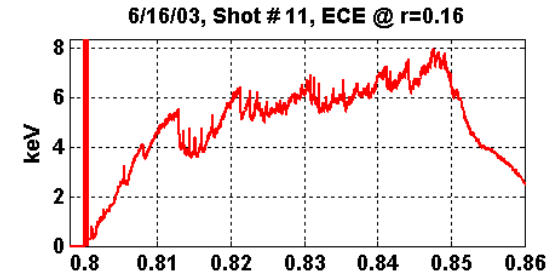
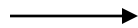
$W \sim 25$ J, $P_{\text{rad}} \sim 17$ kW (40 kW injected)

In Low Density QHS Discharges, 'Crashes' Are Observed in Stored Energy and T_{rad}

Stored
Energy



T_{rad}



Trapped electron modes; electron velocity anisotropy?

Concluding Remarks

- Central T_e and stored energy increase linearly with power, in agreement with $\chi \sim 1/n$ model.
- For constant power, T_e is roughly independent of density, also in accord with $\chi \sim 1/n$ model.
- Model is consistent with H_{α} measurements that show D is roughly independent of power, but depends on $1/n^{0.6}$
- At low density, increases in stored energy are commensurate with energetic trapped population
 - Hard x-ray data
 - Non-thermal ECE emission
 - Outboard resonance mirror returns to proper scaling
- QHS shows higher absorption efficiency and higher X-ray flux than Mirror at low density. At high density, absorbed power falls off at $n > 2 \times 10^{18} \text{ m}^{-3}$.
- Hence, superthermal electrons at low density and degraded absorption at high density account for discrepancy of stored energy with $\chi \sim 1/n$ model.

- Thomson scattering shows a centrally-peaked electron temperature of ~ 600 eV at a line averaged density of $\sim 1.5 \times 10^{18} \text{ m}^{-3}$
- Differences in flows and damping have been observed for thermal plasmas between QHS and Mirror; two timescales
 - QHS slower damping, faster flow, for less drive than mirror
- Superthermal electrons may be drive in stored energy energy drops observed in low density QHS operation

Differences are observed between QHS and Mirror Modes for both thermal and non-thermal plasmas

Near term goals are increasing heating power and plasma density to further understanding of the role and modeling of anomalous transport in quasi-symmetric configurations