



Anomalous Transport Modeling of HSX Plasmas



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Overview

The Weiland ITG/TEM anomalous transport model [1,2] is tested in HSX stellarator plasmas with dominant electron heating

- With geometry simplifications in the quasi-helically symmetric configuration, ITG/TEM predicted stability is comparable to 3D gyrokinetic calculations
- Predicted steady state profiles of electron temperature and density are in moderate agreement with experimental profiles

ITG/TEM Stability

3D ITG/TEM ballooning calculations have demonstrated most unstable eigenmodes in HSX centered about low field, bad curvature region [3,4]

- Dominant particle trapping comes from helical ripple, ε_H (0.14 at $r/a=1$)
- Reduced connection length, $L_c = q_{\text{eff}}R = R/|N-m| \approx 1/3R$
- Low density ECRH plasma \rightarrow Low electron collisionality

$$v_{*e} = \frac{v_e / \varepsilon_H}{\varepsilon_H^{1/2} \frac{v_{Te}}{q_{\text{eff}}R}} \leq 0.1$$

- Max curvature larger than $1/R$

$$\kappa_{N,\text{max}} \sim 1/45 \text{ cm}^{-1} \approx 3 \times 1/R \quad (R=1.2 \text{ m}), \nabla B/B \text{ follows similarly}$$

ITG/TEM Weiland Model

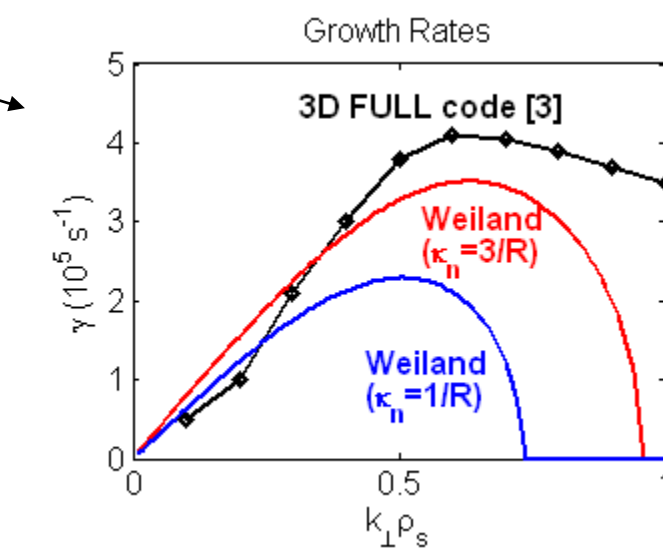
- Fluid toroidal ITG and TEM stability model [1]
- Simplest version includes single ion species, trapped electrons

$$[\chi_e, \chi_i, D, \omega, \gamma] = \frac{\rho_s^2 c_s}{L_n} F\left(\frac{R}{L_{Te}}, \frac{R}{L_{Ti}}, \frac{R}{L_n}, \frac{T_e}{T_i}, f_i, k_{\perp} \rho_s\right)$$

- With geometry approximation, provides reasonable comparison to 3D gyrokinetic FULL code calculations in scaled HSX [3]

Parameters

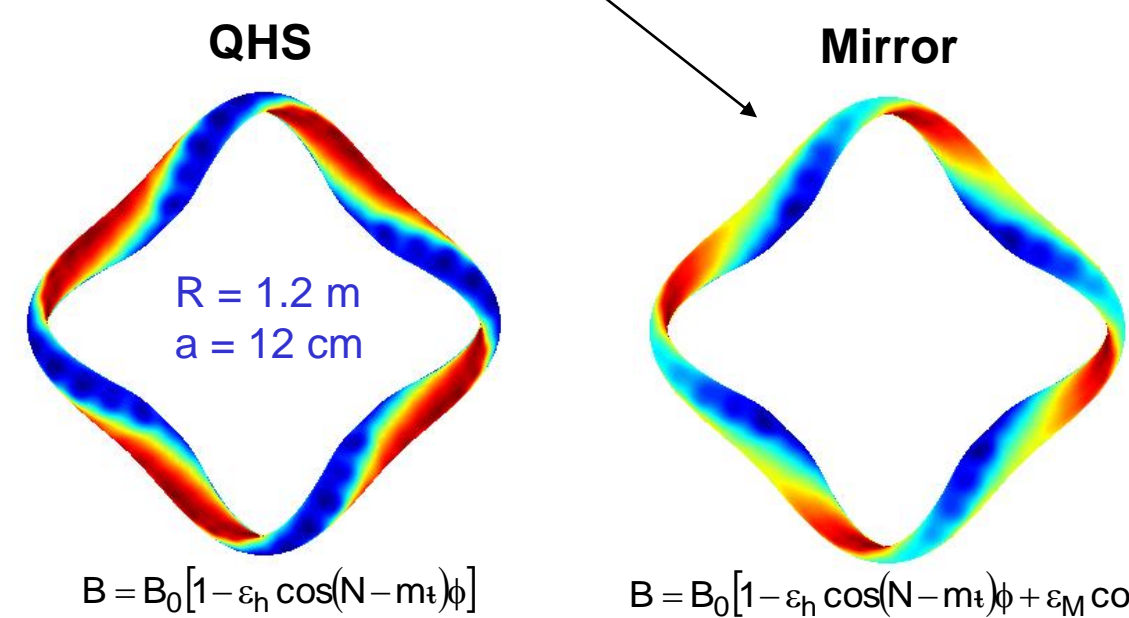
$B = 1.5 \text{ T}$
 $R = 3.6 \text{ m}, a = .36 \text{ m}$
 $r/a = 0.86$
 $n_e = n_i = 1.9 \times 10^{13} \text{ cm}^{-3}$
 $T_e = T_i = 300 \text{ eV}$
 $R/L_n = 13.3$
 $\eta_e = \eta_i = 2.66$
 $Z_{\text{eff}} = 1$
 $\beta = 0, v_e = 0, \omega_{E \cdot B} = 0,$



Experimental Details

HSX has a helical axis of symmetry in $|B|$ and a resulting predicted very low level of neoclassical transport

For experimental flexibility, the quasi-helical symmetry can be broken by adding a mirror field



HSX Plasmas Heated With 28 GHz ECRH

$B = 0.5$ (1.0) Tesla

$P_{\text{ECRH}} = 25\text{-}100 \text{ kW}$, 2nd harmonic X-mode (fundamental O-mode)

$\langle n_e \rangle \leq 3$ (6) $\times 10^{12} \text{ cm}^{-3}$

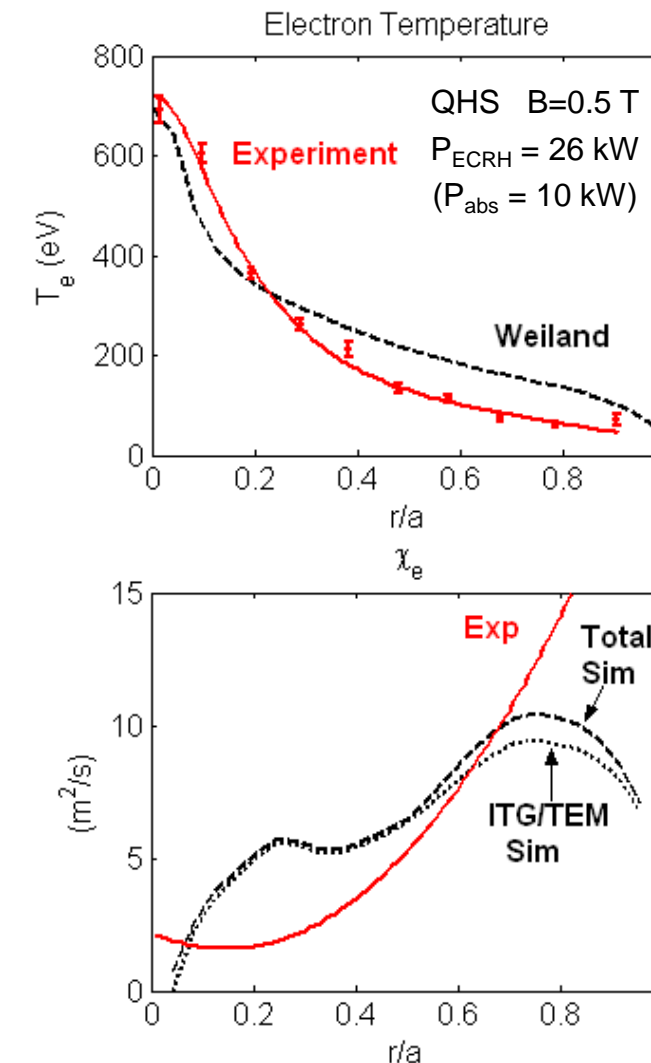
$T_e(0) \sim 700$ (1400) eV $\langle \beta_e \rangle \approx \frac{1}{4000} < \frac{m_e}{m_H} < \beta_e(0) \approx \frac{1}{400}$

$T_i \sim 20\text{-}40 \text{ eV}$ from impurity Doppler spectroscopy

T_e Simulations

$$\frac{3}{2} n_e \frac{\partial T_e}{\partial t} = \frac{1}{r} \frac{\partial}{\partial r} m_e \chi_e \frac{\partial T_e}{\partial r} + P_{\text{ECRH}} - P_{\text{ei}}$$

- Electron energy source from ECRH
 - Total absorbed power from measurements
 - Profile from ray tracing
- Radiation neglected
- $\chi = \chi_{e,\text{NC}} + \chi_{e,\text{an}}$
- Neoclassical contribution from fit to Monte Carlo mono-energetic diffusion coefficients [5]
- Boundary conditions taken from experiment
- Integrated with Matlab PDE solver
- *No free fit parameters have been used*



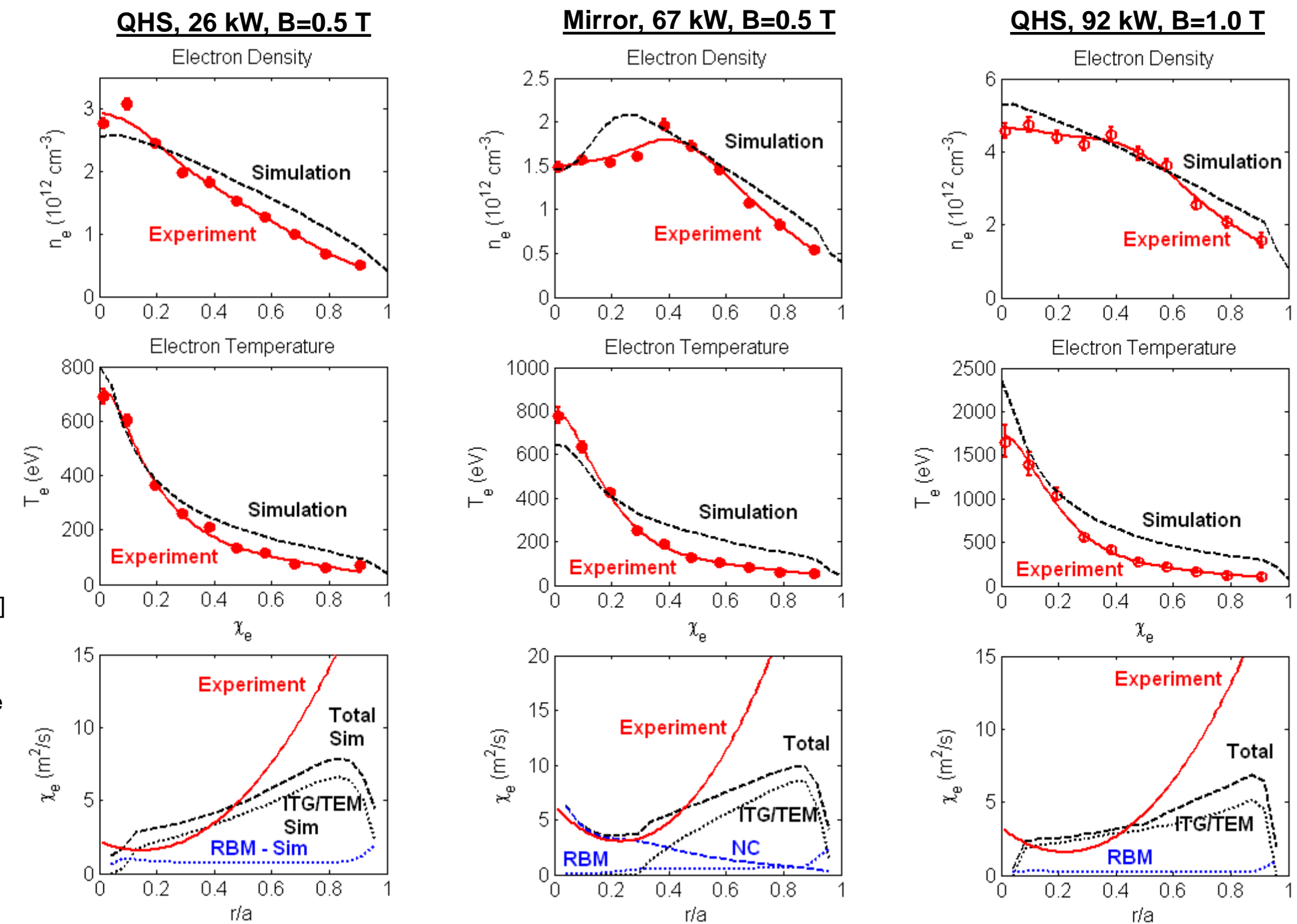
$$\frac{3}{2} n_e \frac{\partial T_e}{\partial t} = \frac{1}{r} \frac{\partial}{\partial r} m_e \chi_e \frac{\partial T_e}{\partial r} + P_{\text{ECRH}} - P_{\text{ei}}$$

$$\frac{3}{2} n_i \frac{\partial T_i}{\partial t} = \frac{1}{r} \frac{\partial}{\partial r} m_i \chi_i \frac{\partial T_i}{\partial r} + P_{\text{ei}} - P_{\text{ex}}$$

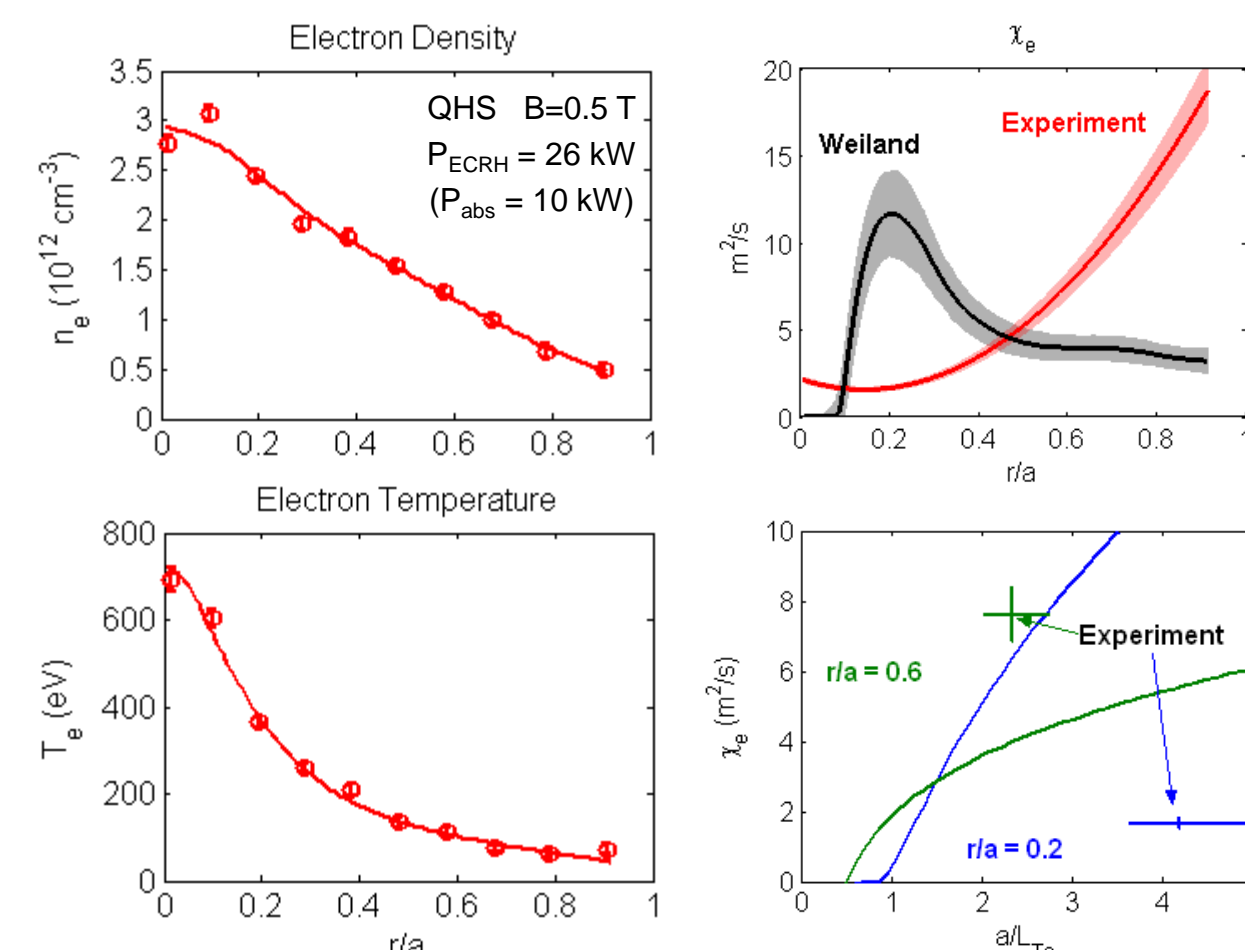
$$\frac{\partial n_e}{\partial t} = \frac{1}{r} \frac{\partial}{\partial r} r D \frac{\partial n_e}{\partial r} + S_{\text{particle}}$$

- Particle source from gas puff and recycling
 - Profile based on 3D neutral gas simulations
 - Magnitude allowed to adjust in simulation to match $\langle n_e \rangle_{\text{sim}}$ to $\langle n_e \rangle_{\text{exp}}$
 - Ion energy sink from i-n charge exchange
 - Updating ambipolar E_r every 50 μs
 - Also including resistive ballooning mode contribution, as used in MMM [2]
 - 3D stability calculations [3] in non-symmetric configuration demonstrate eigenmodes and growth rates that are very similar to the quasi-helically symmetric case
 - Minor differences due to small local differences in $|B|$ & κ in low field, bad curvature region
- \rightarrow Use of same anomalous model and appropriate neoclassical model in simulations

T_e, n_e, T_i Simulations

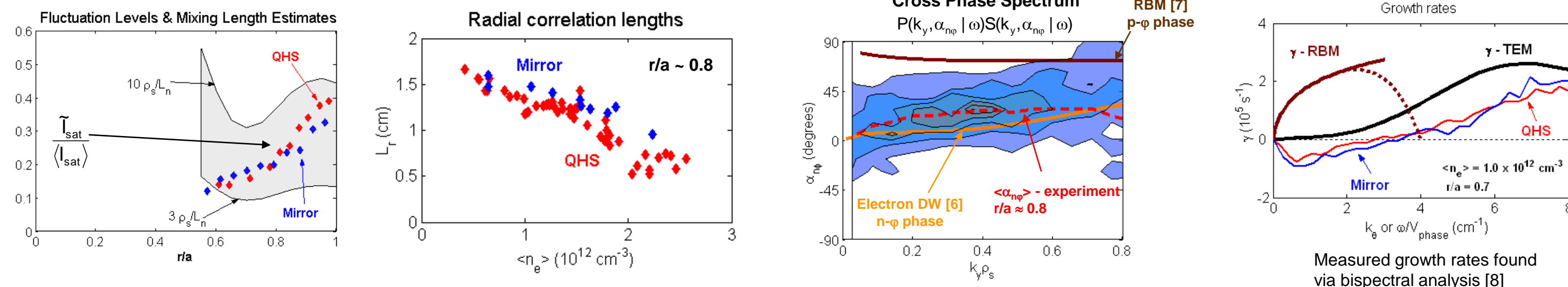


Comparison of Predicted and Experimental χ_e



Turbulence Measurements

Measured Turbulence Characteristics are Very Similar in QHS & Mirror With Same Heating Power (50 kW Injected)



Future Work

- Weiland stability comparison to 3D gyrokinetic code at very limited conditions
- \rightarrow Perform stability calculations using 3D gyrokinetic codes (GS2) under broader conditions
- \rightarrow Calculate quasi-linear fluxes with gyrokinetic code until nonlinear simulations become available

References

- [1] H. Nordman, J. Weiland, & A. Jarmén, Nucl. Fusion **30**, 983 (1990).
- [2] G. Bateman et al., Phys. Plasmas **5**, 1793 (1998).
- [3] T. Rafiq & C.C. Hegna, Phys. Plasmas **12**, (2005).
- [4] G. Rewoldt, L.P. Ku, & W. Tang, Phys. Plasmas **12**, (2005).
- [5] J.N. Talmadge et al., Fus. Sci. & Tech. **46**, 255 (2004).
- [6] W. Horton, Phys. Fluids **19**, 711 (1976).
- [7] J. Callen, UW-CPTC 05-9 (2005).
- [8] J.S. Kim et al., Phys. Plasmas **3**, 3998 (1996).