

Abstract

Second harmonic extraordinary mode ECH is used in the HSX stellarator at 0.5 T to break down and heat the plasma. To measure the absorbed power a set of absolutely calibrated microwave diodes have been installed inside the machine. In the QHS and Mirror configuration, the absorption efficiency is high (about 0.9) and drops (0.6) in the anti-Mirror mode. A comparison with ray tracing predictions is made.







Length **Auxiliary Coils**

Magnetic Field Flattop 0.2 s

• Mod |B| ripples are small in QHS configuration

• In Mirror mode the heating takes place on mod |B| hill while in anti-Mirror mode in the well.

EC Heating in HSX



• Microwave power (up to 100 kW at 28 GHz) is used for neutral gas breakdown followed by heating of plasma at the second harmonic of electron cyclotron frequency.

Toroidal Angle φ

Linear polarized beam with **E** perpendicular to **B** is launched from the low magnetic field side and is focused on the plasma center in a spot of 4 cm in diameter.

Ray-Tracing Equations

$\frac{d\vec{r}}{dt} =$	$= -\frac{\partial G}{\partial \vec{k}}$	<u>Cold plasma dispersion relation</u> $G(\vec{r}, \vec{k}, \omega) = \frac{c^2 \vec{k}^2}{\omega^2} - N^2(\vec{r}, \vec{k}, \omega) = 0$
$\frac{d\vec{k}}{dt}$	$=\frac{\partial G}{\partial \vec{r}}$	$\frac{\Delta M^{2}}{\frac{Altar-Appleton-Hartree formula}{\omega_{pe}^{2}(\omega^{2}-\omega_{pe}^{2})/\omega^{2}}}$
$\frac{dP}{dt} =$	$-2 \cdot \operatorname{Im} k_j \cdot V_{gr}$	$\omega^{2} - \omega_{pe}^{2} - 0.5 \cdot \Omega_{e}^{2} \sin^{2} \theta \mp \Delta$ $ \cdot \cos(\alpha) \cdot P \qquad \Delta = \sqrt{0.25 \cdot \Omega_{e}^{4} \sin^{4} \theta + \Omega_{e}^{2} \omega^{-2} (\omega^{2} - \omega_{pe}^{2})^{2} \cos^{2} \theta}$

 Im/k_i - absorption coefficient; V_{gr} - group velocity; α - angle between V_{gr} and k_i .

Electron Cyclotron Damping

• Resonance condition: $\omega = n \cdot \Omega_e + k_{\parallel} \cdot v_{\parallel}$

where $\Omega_e = \frac{eB}{m_o c_v 1 + \left(\frac{v_{th}}{c_o}\right)^2}$

• Extraordinary wave at the second harmonic of Ω_{e} $\frac{\text{Im} |k_e|}{k_e} = \frac{2\pi}{15} q \cdot \left(\frac{3 - 0.5q}{3 - q}\right)^2 \cdot z^{\frac{5}{2}} \cdot e^{-z} \qquad \text{for} \qquad \theta = \frac{\pi}{2}$

 $\frac{\mathrm{Im} \mid k_e \mid}{k_e} = \frac{\pi}{4} q \cdot \left(\frac{3 - 0.5q}{3 - q}\right)^2 \cdot \left(\sqrt{1 + 2z/a}\right)^{\frac{3}{2}} \cdot I_{\frac{3}{2}} \left(a\sqrt{1 + 2z/a}\right) \cdot e^{-(z+a)} \qquad \qquad \left|\frac{\pi}{2} - \theta\right| << 1$

where $a = \frac{mc^2}{T_e} \cdot \frac{3 - 2q + 0.5q}{3 - q}$ $q = \frac{\omega_{pe}^2}{\Omega_e^2}$ $z = \frac{2\Omega_e - \omega}{\Omega v_e^2/c^2}$





data





Absorption of X-Wave at the Second Harmonic in HSX

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Perpendicular and Oblique Propagation

• At $\Theta = 90$ degs. absorption occurs at $\Omega_{\rm e} > \omega$ • Absorption efficiency

- drops with **k**_{parallel}
- At large **k**_{parallel} the absorption is symmetric with respect to resonance point



Temperature effect

- absorption shape is broadened with temperature
- At $\Theta = 90$ degs. the maximum is moved inward due to relativistic mass increase

These profiles

are used in the

ray tracing

calculations



• Plasma density profile is interpolated interferometer

Single-pass Absorption

Magnetic Field Scan



• Shift out of the central resonance reduces the absorption

Reflection From The Wall







- Absorbed power profile is not broaden much
- Energy per particle in the center rises



MD #6/ $B_2^{-} - B_3^{-}$

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Absorbed Power Profile



• Each antenna consists of an open K-band waveguide, attenuator and microwave detector

Toroidal Angle, degs.	Distance from ECH port, m	Area Of HSX Helical Cut, m ²
6.3	0.181	0.12
36.0	0.942	0.1188
69.3	1.675	0.1297
-69.3	1.675	0.1297
103.5	2.642	0.1264
-103.5	2.642	0.1264

Calibration of Diodes



• The microwave source and the line are calibrated with thermistor mount, first, and then the diodes.

Calibration Curves



- Six diodes are almost identical
- up to about $30 \,\mu W$ of input power

Power at each toroidal location

Atten	uator	Amplifier	• Power at location:
Quartz Window	μw De	tector	$P_{non} = \frac{U_i}{G_i} \cdot S$

 U_i – acquired signal; G_i – gain of the amplifier; S_i – diode sensitivity; Att – attenuation coefficient of the attenuator; trs_w – transparency coefficient of the quartz window; Wg – area of waveguide cross-section; S_{v} – area of the HSX helical cut



• Antennas are mounted on 5" extension nipples



Raw signals



In shot #36: $n_e = 1.5 * 10^{18} \text{ m}^{-3}$ $W_E = 22 J$ <u>In shot #37:</u> $n_e = 3.1 \times 10^{18} \text{ m}^{-3}$ $W_{\rm E} = 1.6 \, {\rm J}$

The square law is





In shot #36: $n_e = 1.5 * 10^{18} \text{ m}^{-3}$ $W_{\rm E} = 22 \, {\rm J}$ In shot #37: $n_e = 3.1 * 10^{18} \text{ m}^{-3}$ $W_{\rm F} = 1.6 \, \rm J$

0 0.79 0.8 0.81 0.82 0.83 0.84 0.85 2/18/03, shot #36

0.79 0.8 0.81 0.82 0.83 0.84 0.8

2/18/03. shot #30

Absorption coefficients

 $W_E = 22 J$

In shot #36:

 $n_{e} = 1.5 * 10^{18} \text{ m}^{-3}$

MD #4

0.79 0.8 0.81 0.82 0.83 0.84 0.85

2/18/03, shot #36

0.81 0.82 0.83 0.84 0.8

2/18/03. shot #3

• Absorption coefficient: $\eta = 1 - P_{non}/P_o$, where $P_{non} - P_o$ non-absorbed power, P_0 - power in "cold" plasma discharge (at high plasma density)

Plasma Density Scan

Non-absorbed ECH Power Mirror mode 3/17/03 mode 2/18/03 kW 🗕 🔷 MD #1 ◆ MD #1 • MD #2 • MD #2 • • • × MD #4 → MD #5 <u>∧</u> MD #6 × × Line average densitv. 10⁻¹⁸ m⁻³ Line average density, 10⁻¹⁸ m⁻³

• Launched Power is about 40 kW in both cases

Plasma Density Scan (cont.)

Multi-pass Absorption Efficiency



• In QHS and Mirror configurations the absorption is high (0.8-0.9). At low plasma densities the error of this method can be big enough due to wave interference.

ECE measurements



• In QHS mode the ECE signal is about a factor of 3 higher at low plasma density than at high plasma density. • It might be due to high plasma temperature and some population of superthemal electrons

each toroidal

 $Wg \cdot Att \cdot trs_{w}$